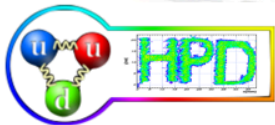




Geometrical scaling from RHIC to LHC energies



Amelia LINDNER

Coordinator: Prof. Dr. Mihai PETROVICI



HADRON PHYSICS DEPARTMENT

**ISAB Meeting, Bucharest, November 1st, 2018
Young Scientist Forum**

Outline

➤ Physics motivation

➤ Geometrical scaling

- $\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$
- The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$
- $\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$

▪ Au+Au & Cu+Cu - RHIC, Pb+Pb & Xe+Xe – LHC

Energy domain: RHIC (BES (7.7 – 39 GeV) + 62.4, 130 & 200 GeV)
LHC (2.76, 5.02 & 5.44 TeV)

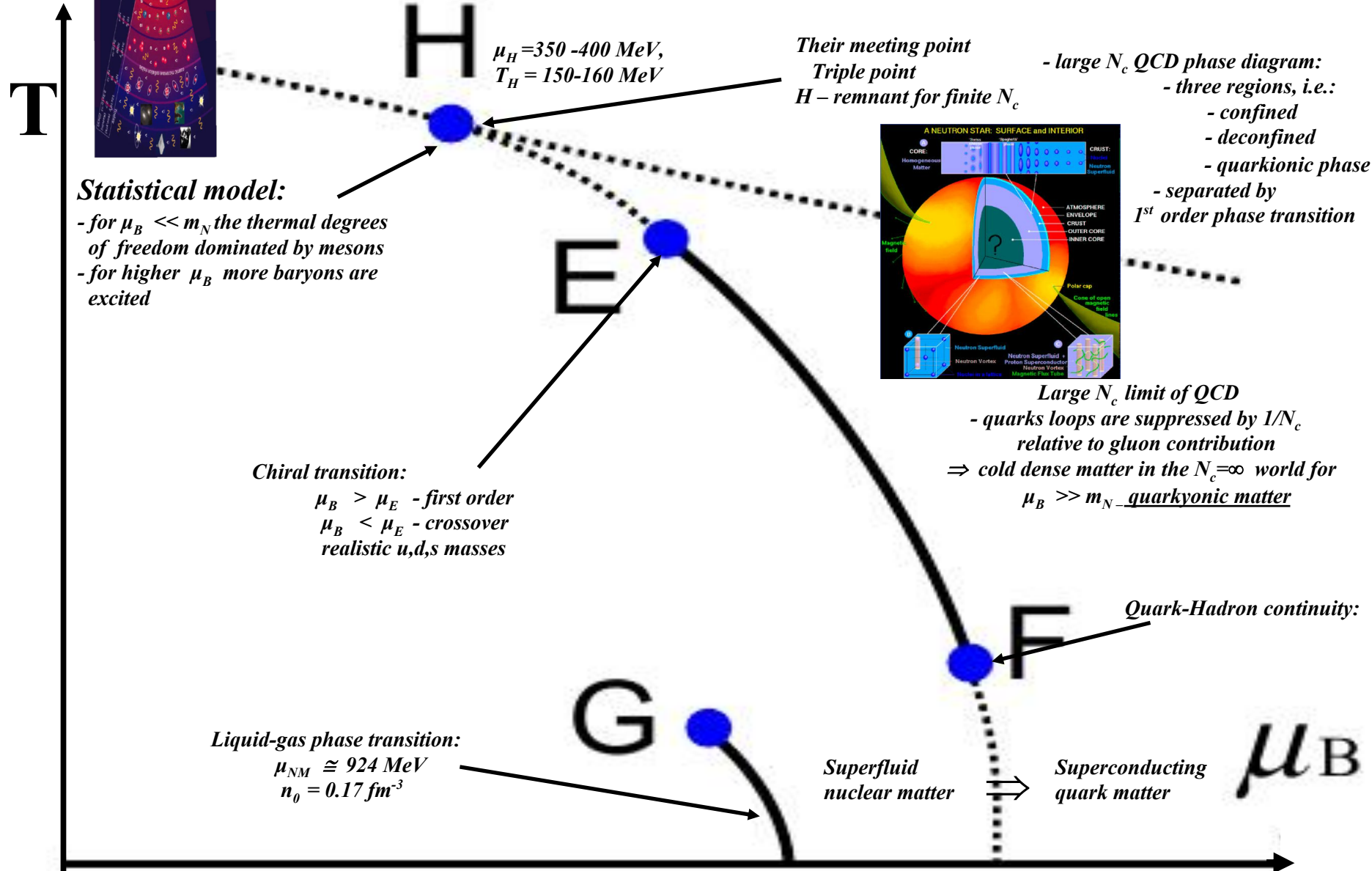
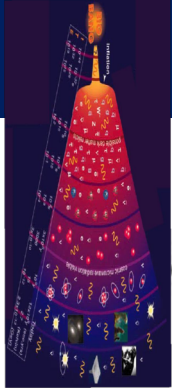
▪ p+p vs. Pb+Pb @ LHC

Energy domain: Pb+Pb (5.02 TeV); p+p (7 TeV)

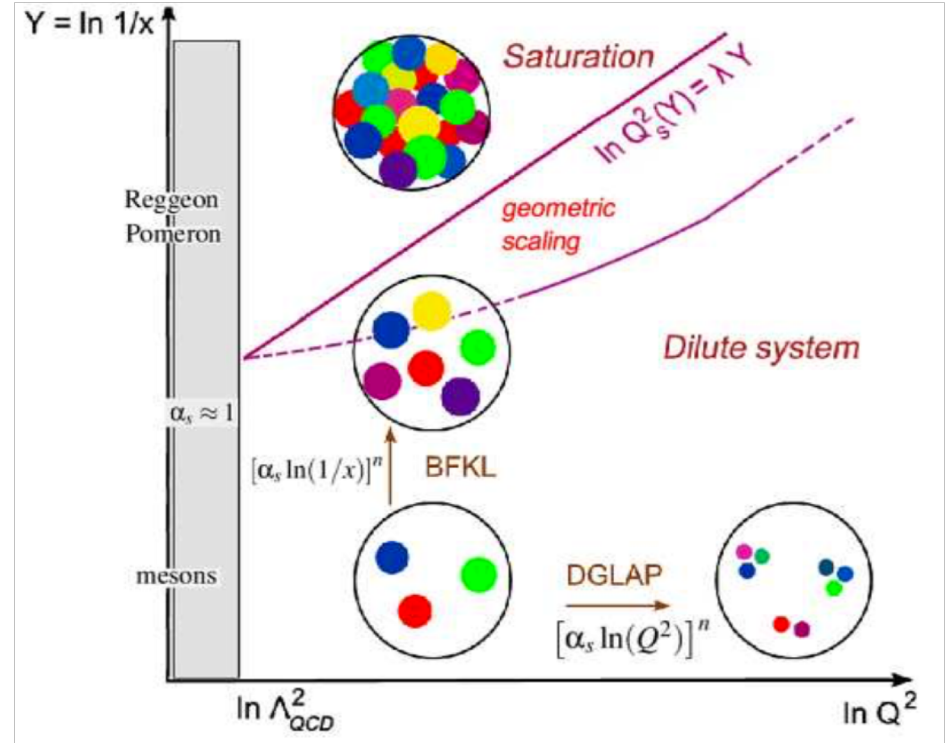
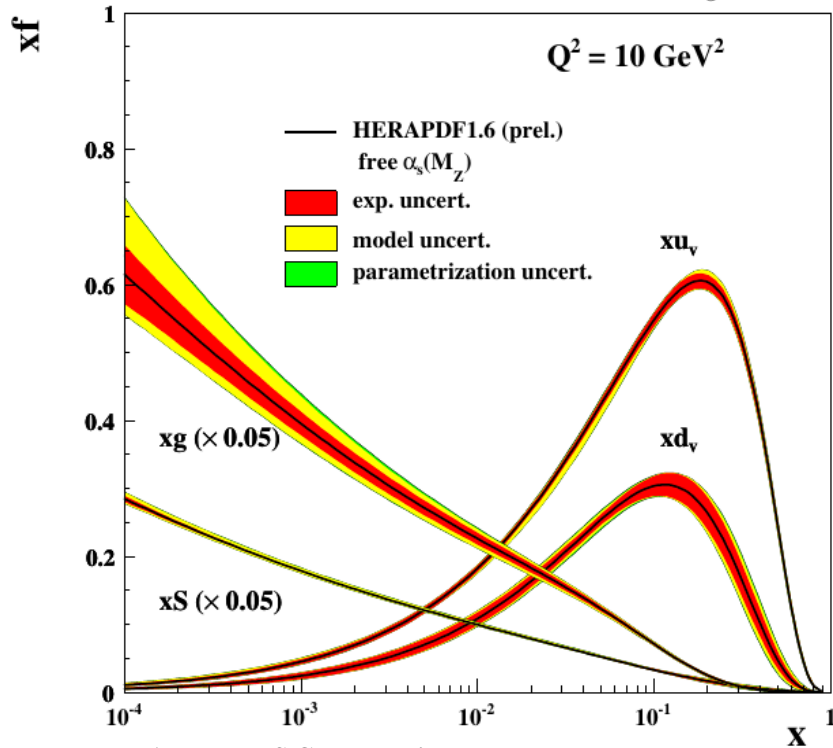
➤ Preliminary Results

➤ Conclusions & Perspectives

Physics motivation



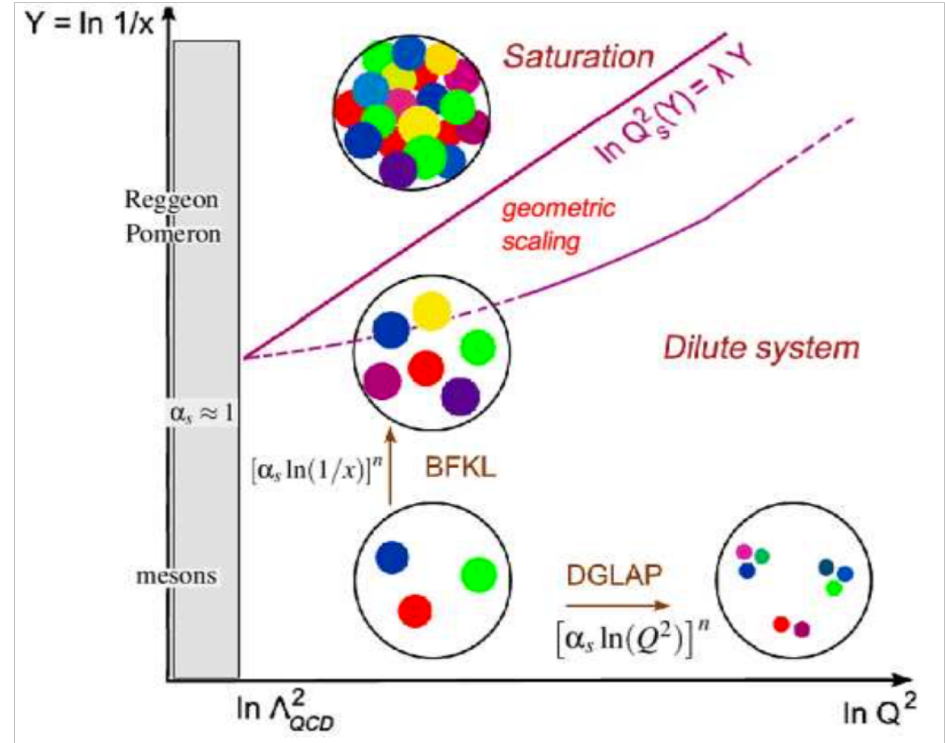
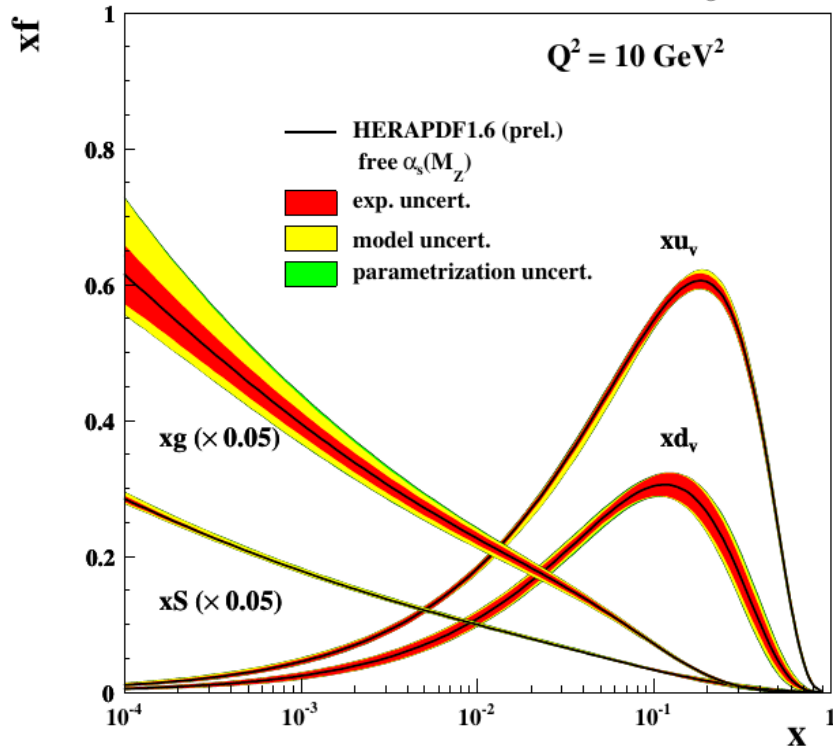
Physics motivation



R.Placakyte, H1 and ZEUS Collaborations, XXXI
PHYSICS IN COLLISIONS, Vancouver, BC
Canada, Aug.28-Sept.1, 2011

D. d'Enterria, Eur.Phys.J. A31(2007)816)

Physics motivation



R.Placakyte, H1 and ZEUS Collaborations, XXXI
PHYSICS IN COLLISIONS, Vancouver, BC
Canada, Aug.28-Sept.1, 2011

D. d'Enterria, Eur.Phys.J. A31(2007)816

Geometrical scaling

Local parton-hadron duality picture
and dimensionality argument

$$\langle p_T \rangle / \sqrt{\frac{dN}{dy} / S_{\perp}} \sim \frac{1}{n\sqrt{n}}$$

$n = \text{no. of charged hadrons produced from a gluon fragmentation}$

$$\langle p_T \rangle / \sqrt{\frac{dN}{dy} / S_{\perp}}$$

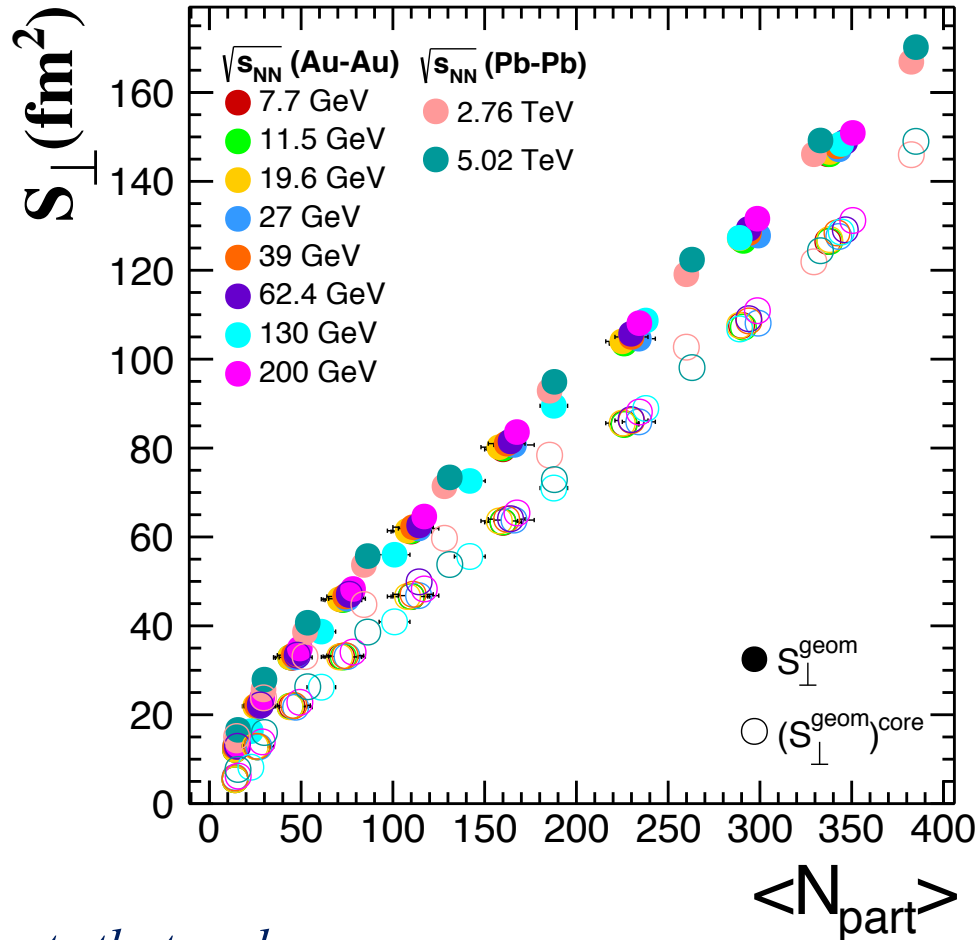
decreases as a function of collision energy and centrality 5

Y.L.Dokshitzer, V.A.Khoze and S.Troian, J.Phys.G 17 (1991) 1585
T. Lappi, Eur.Phys.J. C71 (2011) 1699
E. Levin and A.H. Rezaeian, Phys.Rev.D 83 (2011) 114001

S_{\perp} estimates

Scaling variable: $\sqrt{[(dN/dy)/ S_{\perp}]}$

Glauber Monte Carlo approach



core = the participants that undergo more than a single collision

dN/dy estimates

Scaling variable: $\sqrt{[(dN/dy)/S_{\perp}]}$

BES

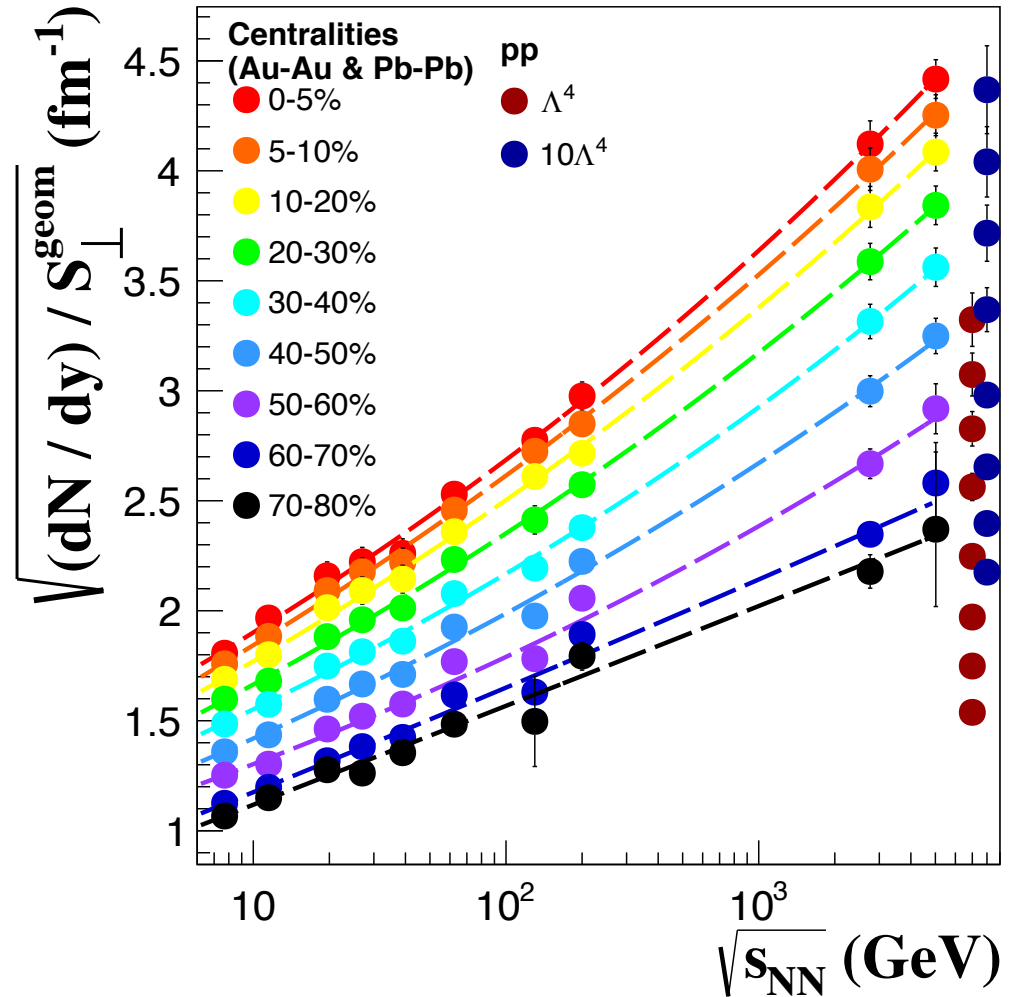
$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+\pi^-)}}{dy} + 2 \frac{dN^{(K^+K^-, p+\bar{p}, \Xi^-+\bar{\Xi}^+)}}{dy} + \frac{dN^{(\Lambda+\bar{\Lambda})}}{dy}$$

RHIC (62.4; 130 & 200 GeV)

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+\pi^-)}}{dy} + 2 \frac{dN^{(K^+K^-, p+\bar{p}, \Xi^-+\bar{\Xi}^+)}}{dy} + \frac{dN^{(\Lambda+\bar{\Lambda}, \Omega^-+\bar{\Omega}^+)}}{dy}$$

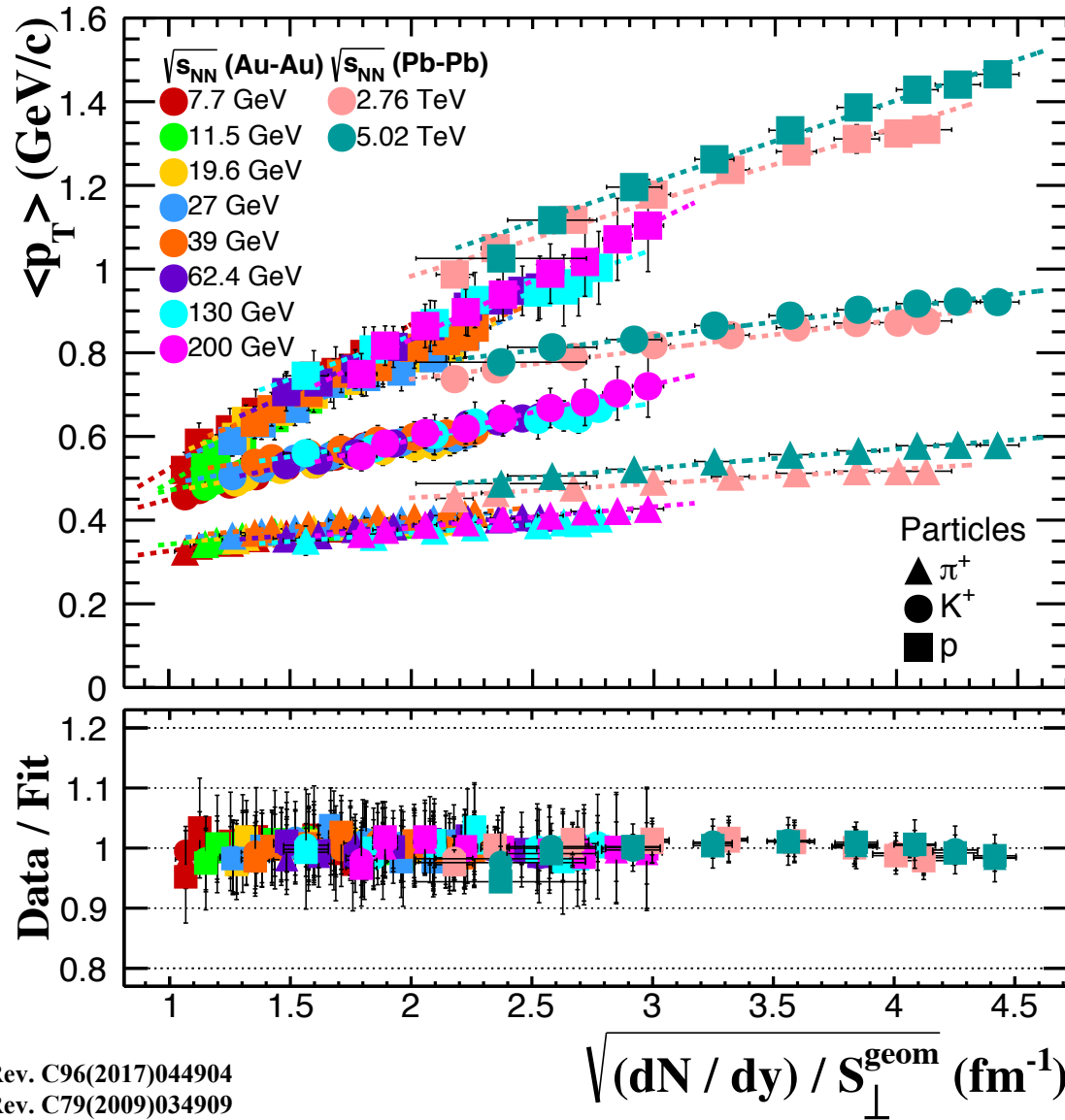
LHC

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+\pi^-)}}{dy} + 2 \frac{dN^{(p+\bar{p}, \Xi^-+\bar{\Xi}^+)}}{dy} + \frac{dN^{(K^+K^-, K_S^0+\bar{K}_S^0, \Lambda+\bar{\Lambda}, \Omega^-+\bar{\Omega}^+)}}{dy}$$



$$\sqrt{(dN/dy)/S_{\perp}^{geom}} = a + b \cdot (\sqrt{s_{NN}})^{2c}$$

$\langle p_T \rangle$ vs. $\sqrt{[(dN/dy) / S_{\perp}]}$



STAR Collaboration, Phys.Rev. C96(2017)044904

STAR Collaboration, Phys.Rev. C79(2009)034909

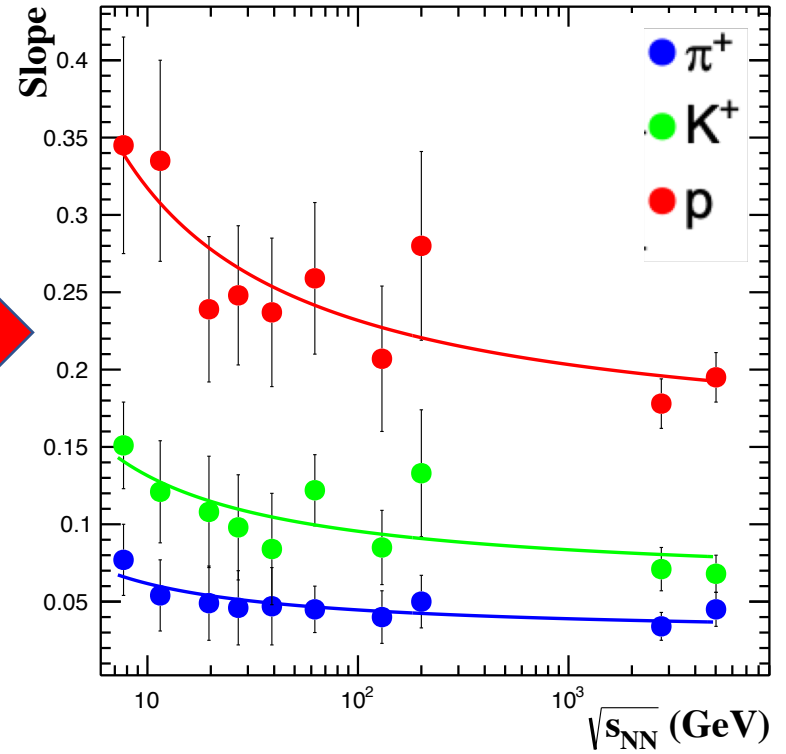
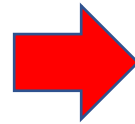
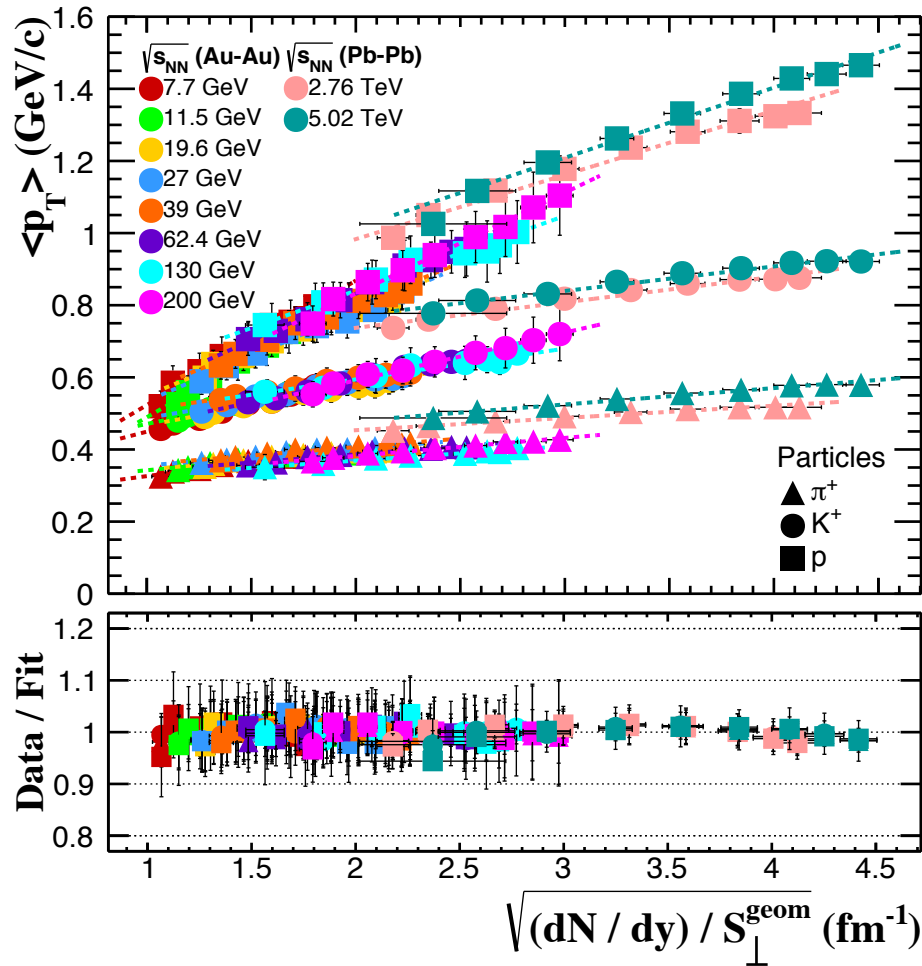
ALICE Collaboration, Phys.Rev. C{88}{2013}{044910}

ALICE Collaboration, Phys.Rev.Lett. 116(2016)222302

ALICE Collaboration, Eur.Phys.J. C75(2015)226

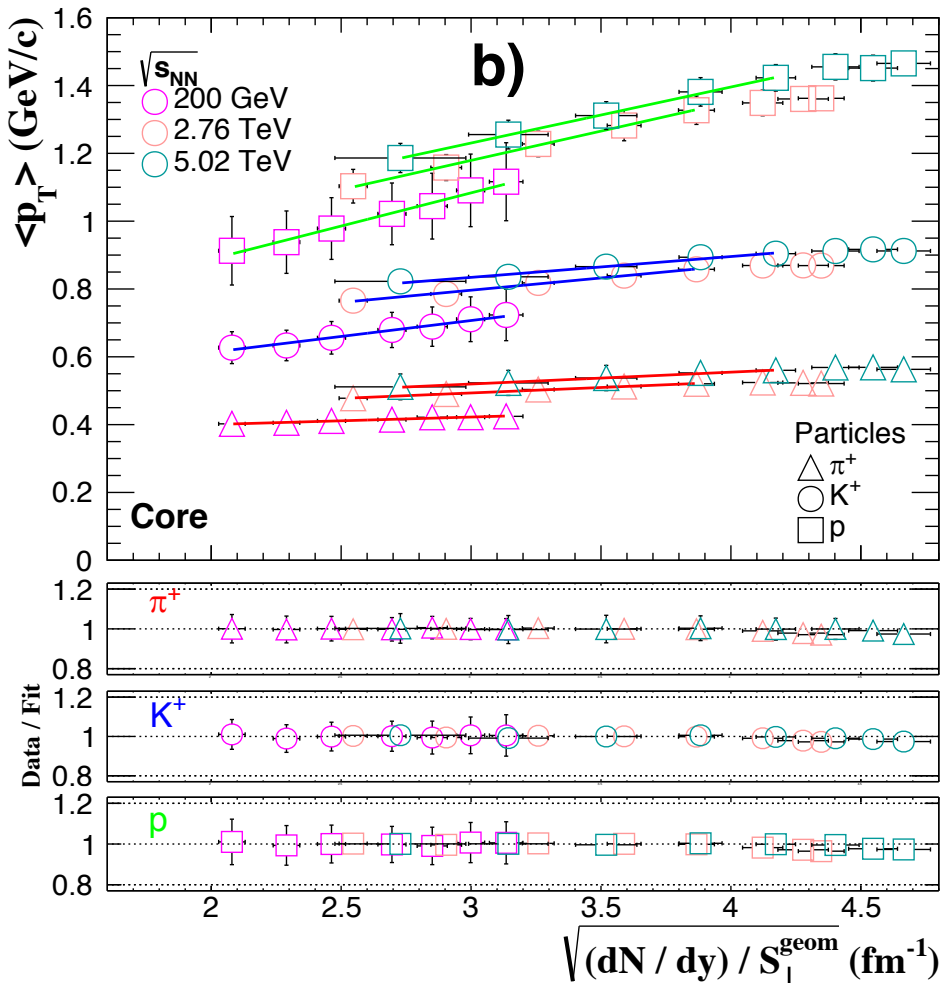
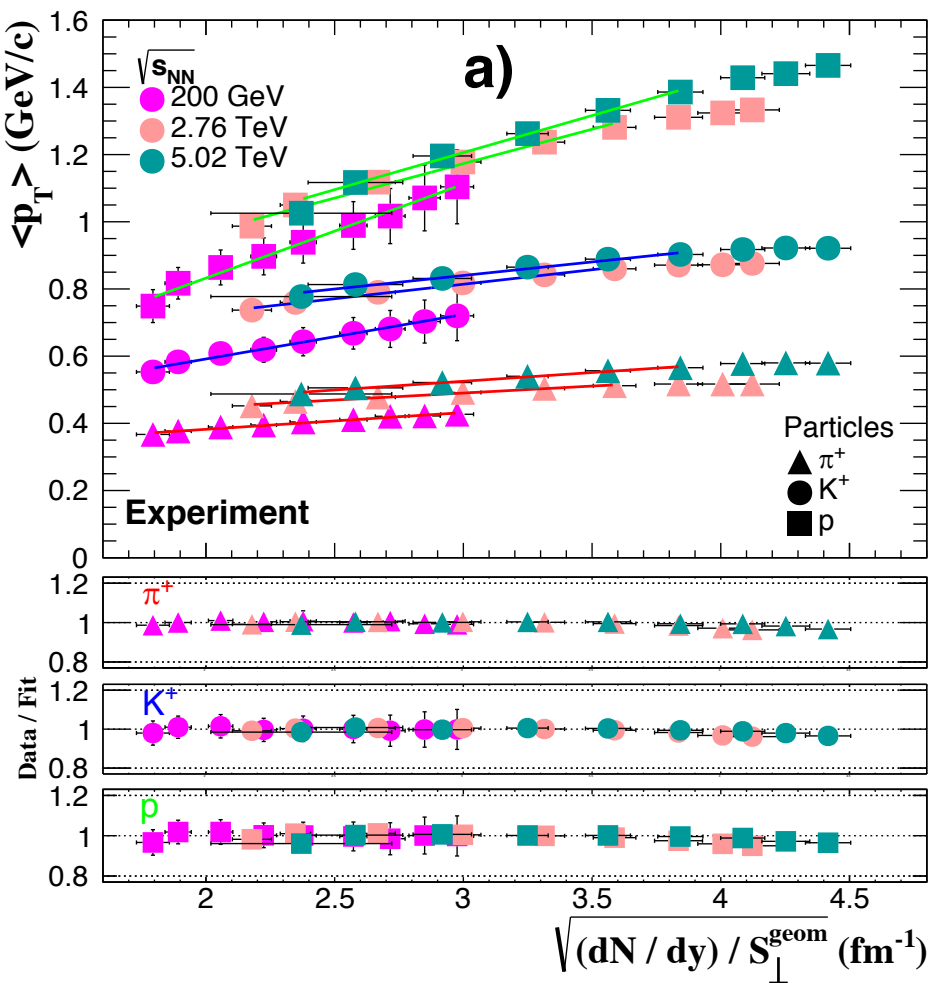
A.K.Dash, ALICE Collaboration, 9th Int. Workshop on MPI at LHC, Dec. 11-15, 2017

$\langle p_T \rangle$ vs. $\sqrt{[(dN/dy)/S_{\perp}]}$

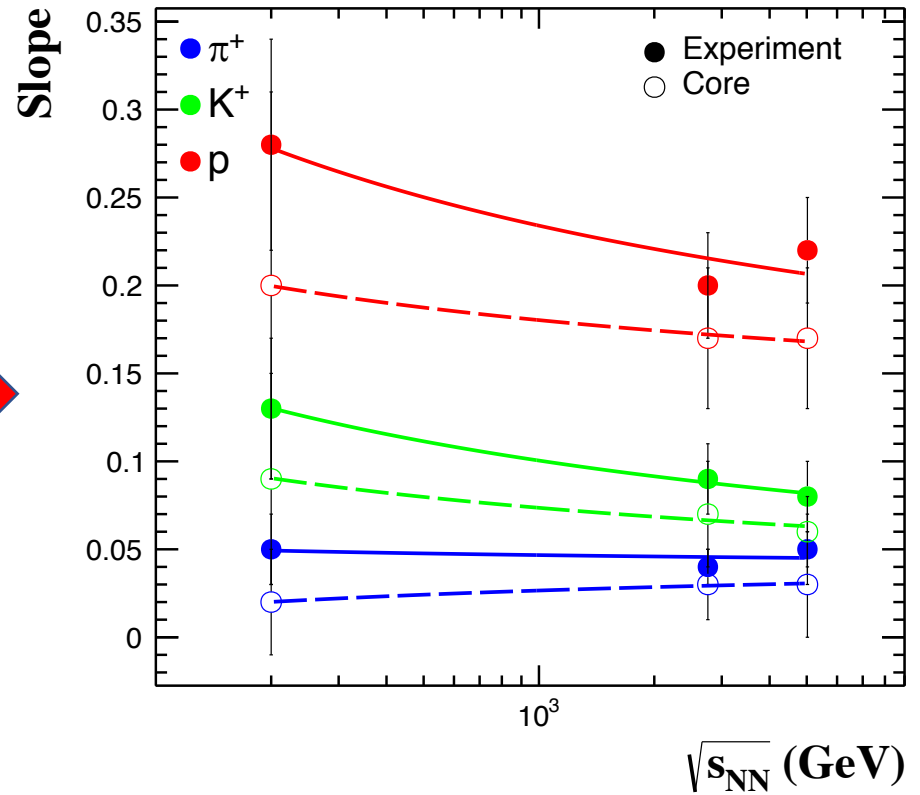
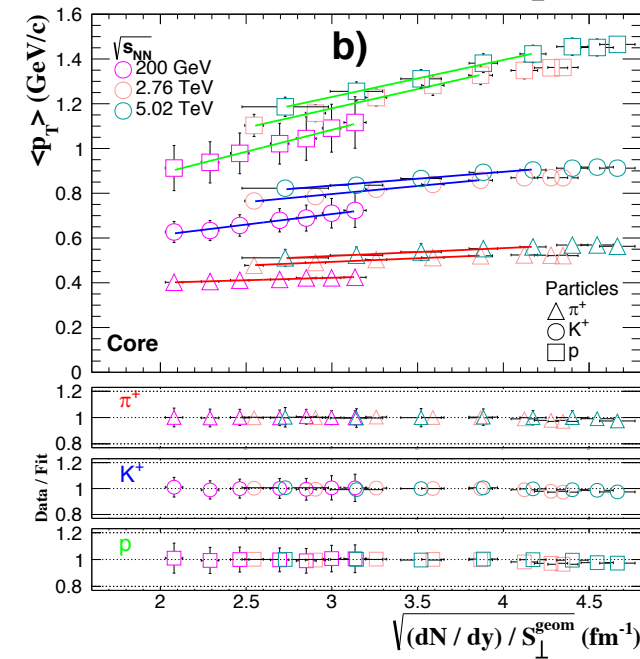
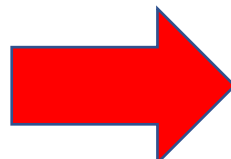
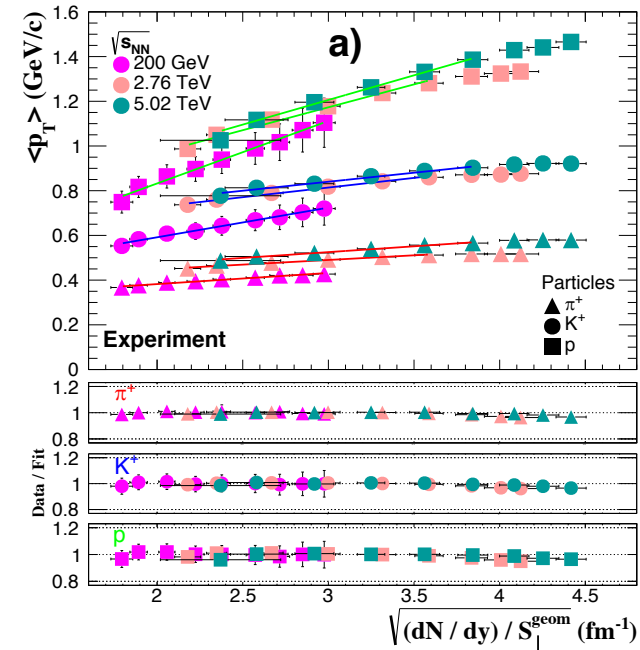


STAR Collaboration, Phys.Rev. C96(2017)044904
 STAR Collaboration, Phys.Rev. C79(2009)034909
 ALICE Collaboration, Phys.Rev. C{88}{2013}{044910}
 ALICE Collaboration, Phys.Rev.Lett. 116(2016)222302
 ALICE Collaboration, Eur.Phys.J. C75(2015)226
 A.K.Dash, ALICE Collaboration, 9th Int. Workshop on MPI at LHC, Dec. 11-15, 2017

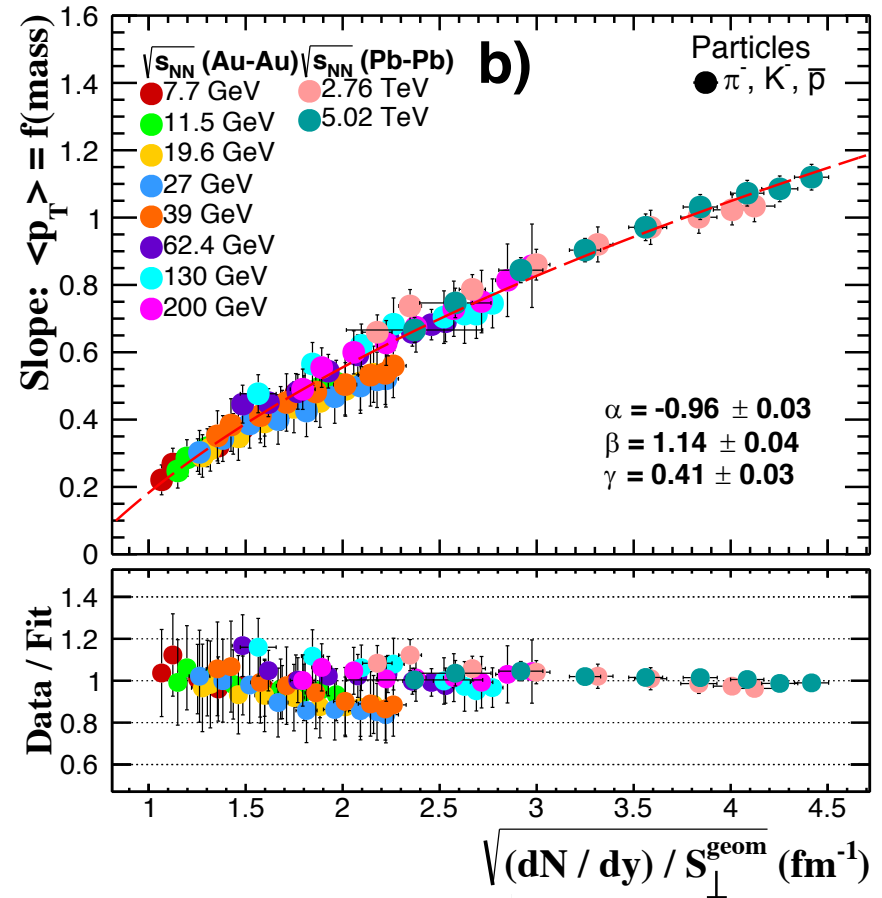
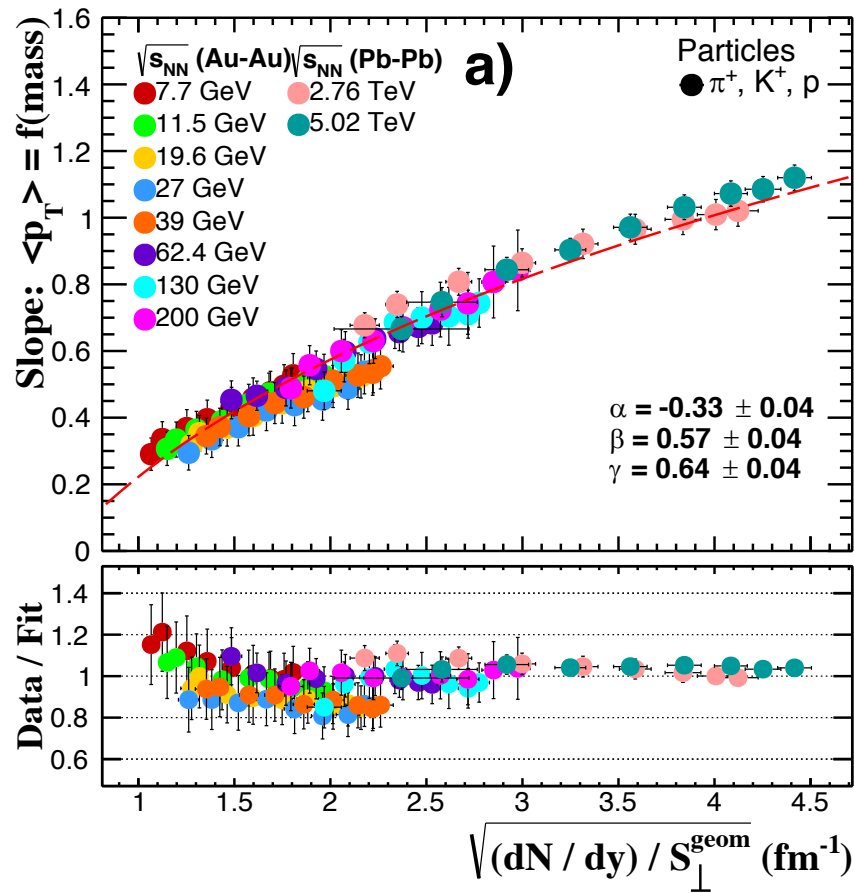
The Core-Corona effect



The Core-Corona effect

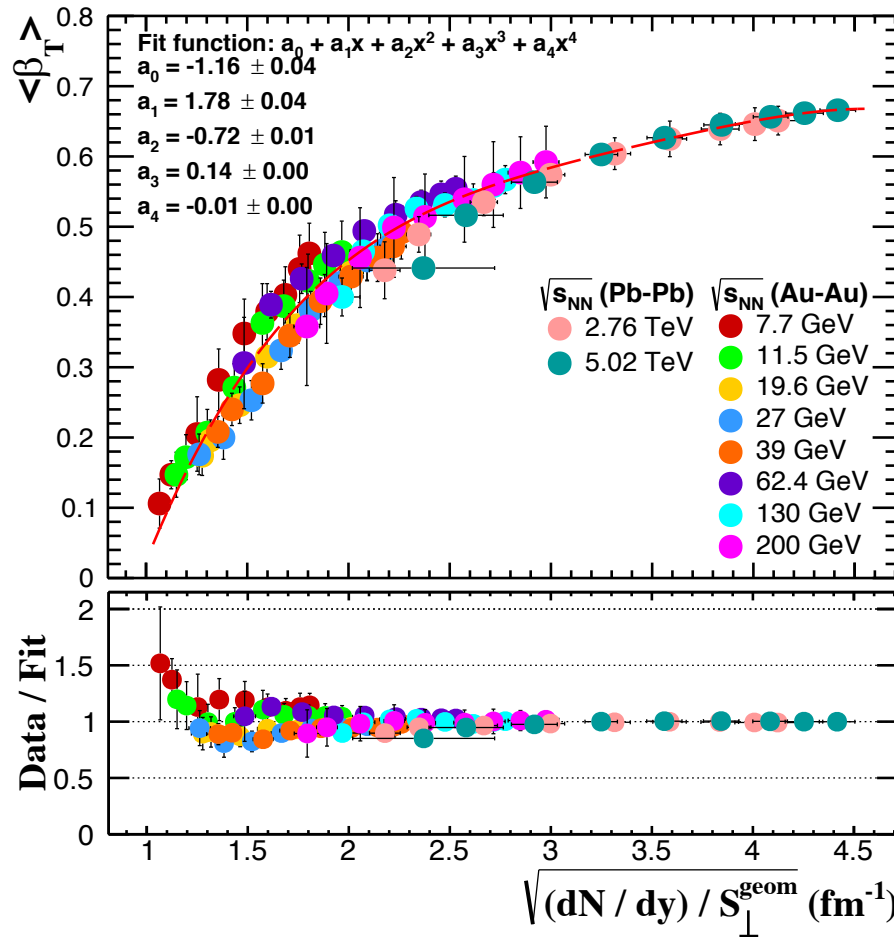


The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $\sqrt{[(dN/dy) / S_{\perp}]}$



$$\text{Slope}_{\langle p_T \rangle = f(\text{mass})} = \alpha + \beta \left(\sqrt{\frac{dN}{dy} / S_{\perp}^{\text{geom}}} \right)^{\gamma}$$

$\langle \beta_T \rangle$ from BGBW fits vs. $\sqrt{[(dN/dy)/S_{\perp}]}$



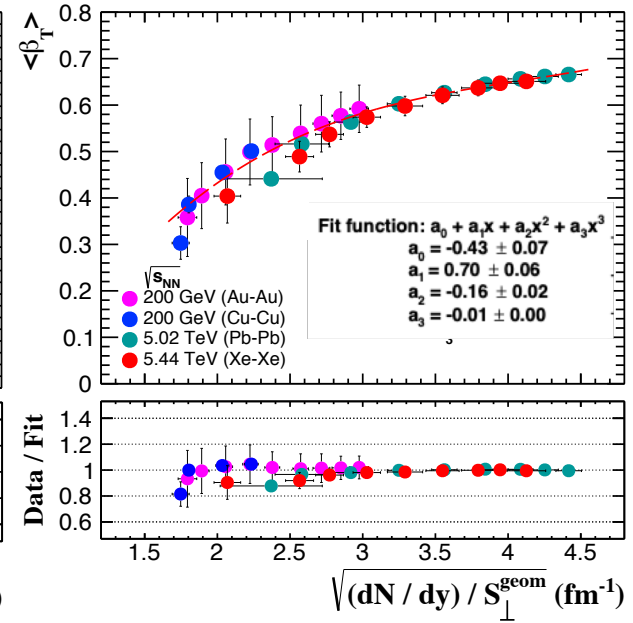
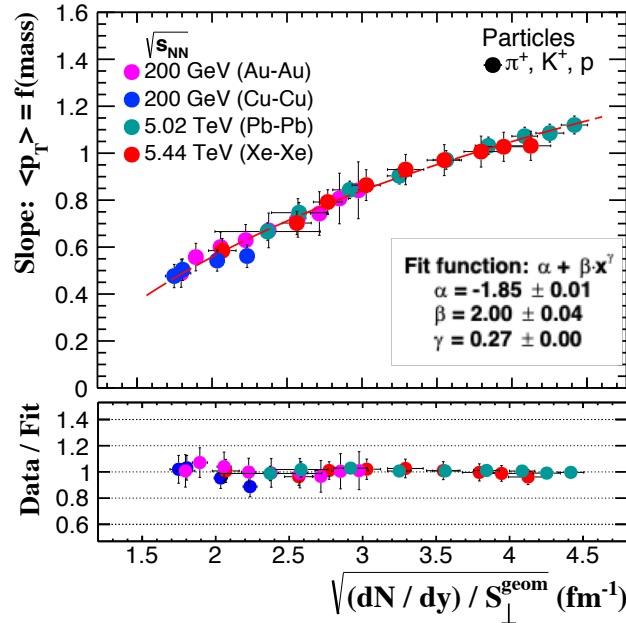
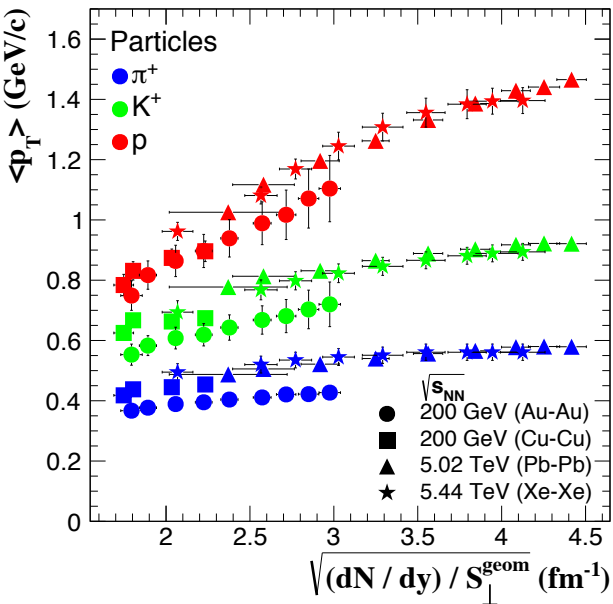
Boltzmann-Gibbs
Blast Wave

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dy dp_T} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}} \right)$$

$$\rho = \tanh^{-1} \beta_T = \tanh^{-1} \left[\left(\frac{r}{R} \right)^n \beta_s \right]$$

Xe+Xe vs. Pb+Pb @ LHC and Cu+Cu vs. Au-Au @ RHIC

Latest Results



Carpathian Summer School of Physics (2018)
 --> AIP - Conference Proceedings

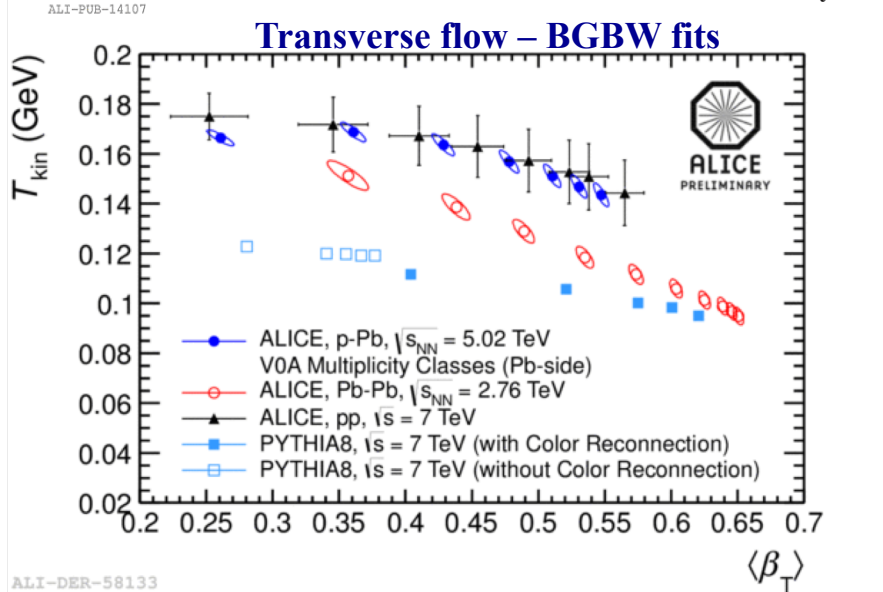
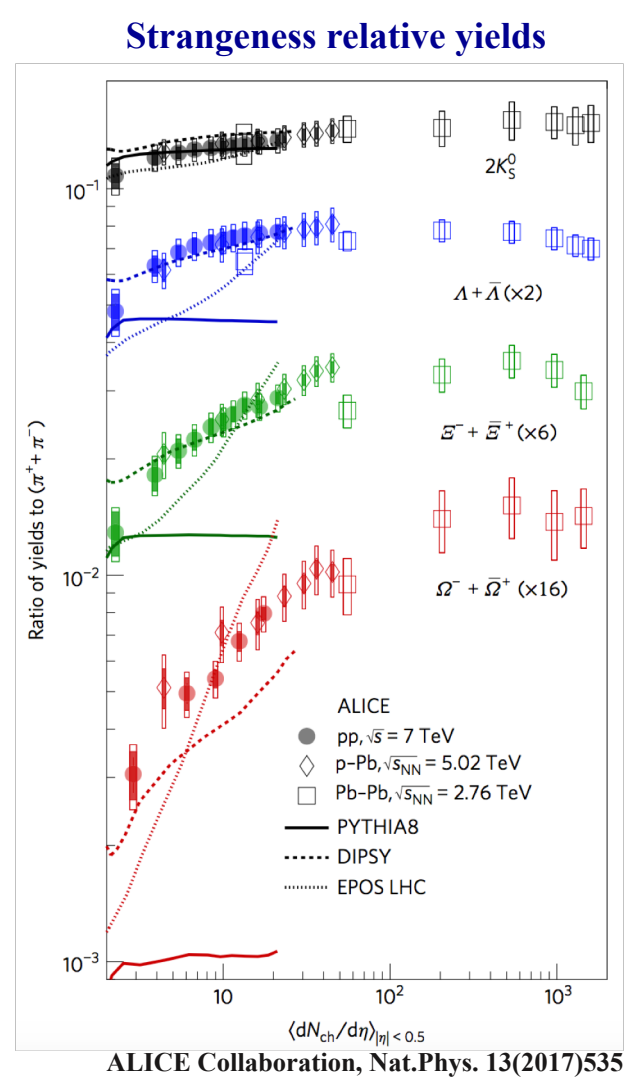
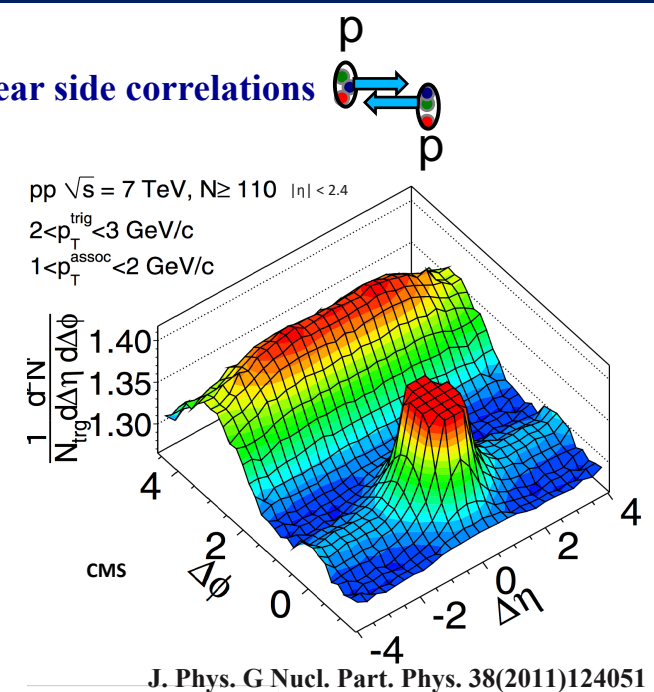
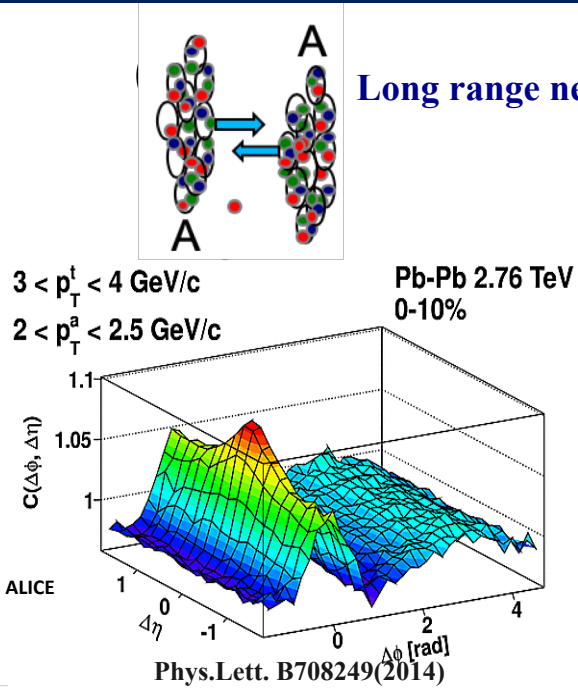
F. Bellini, ALICE Collaboration, Quark Matter 2018

D. S. D Albuquerque, ALICE Collaboration, Quark Matter 2018 STAR Collaboration, Phys.Rev.Lett. 108 (2012) 072301

BRAHMS Collaboration, Phys.Rev.C 94 (2016) 014907

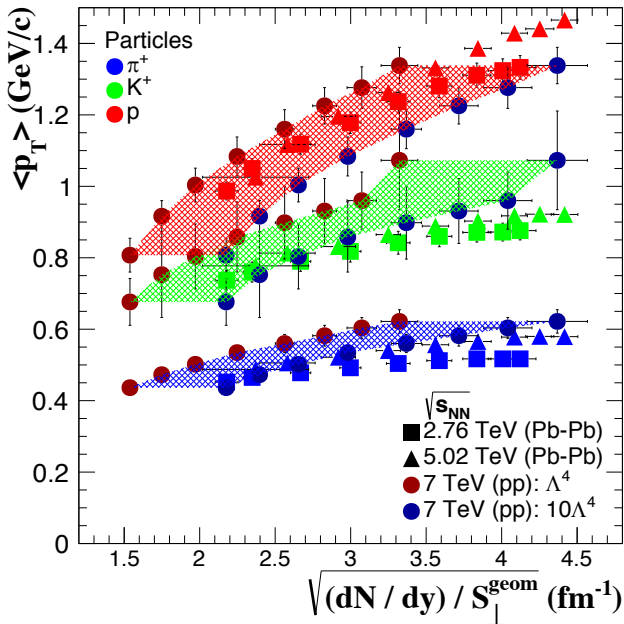
STAR Collaboration, Phys.Rev.Lett. 108 (2012) 072301

p+p vs. Pb+Pb @ LHC

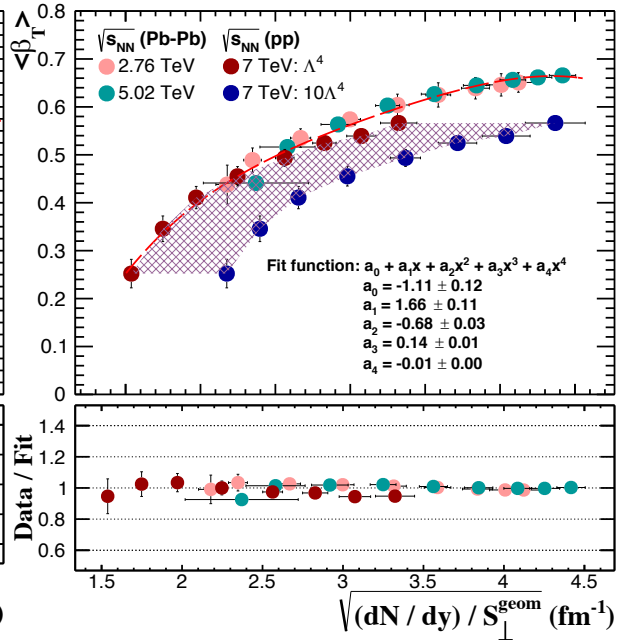
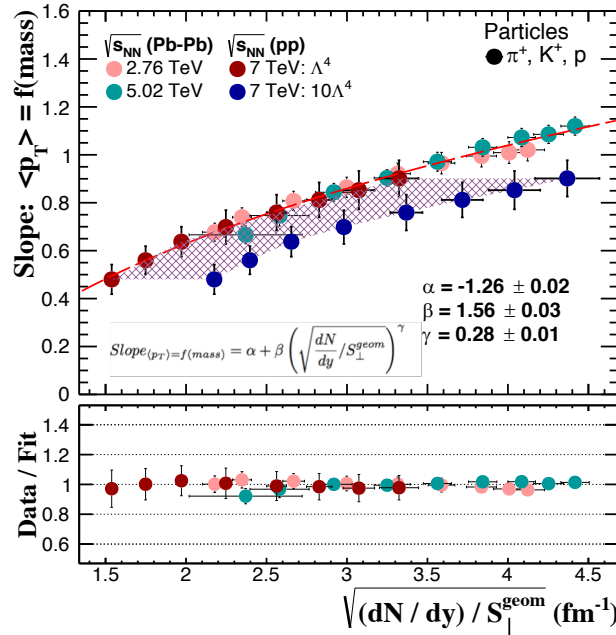


p+p vs. Pb+Pb @ LHC

$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$ $\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



$$S_{\perp}^{pp} = \pi R_{pp}^2 \quad \longrightarrow \quad \varepsilon = \alpha \Lambda_{QCD}^4$$

$\alpha=1$

$$f_{pp} = \begin{cases} 0.387 + 0.0335x + 0.274x^2 - 0.0542x^3 & \text{if } x < 3.4 \\ 1.538 & \text{if } x \geq 3.4 \end{cases}$$

$$x = (dN_g/dy)^{1/3}$$

$$dN_g/dy \approx dN/dy$$

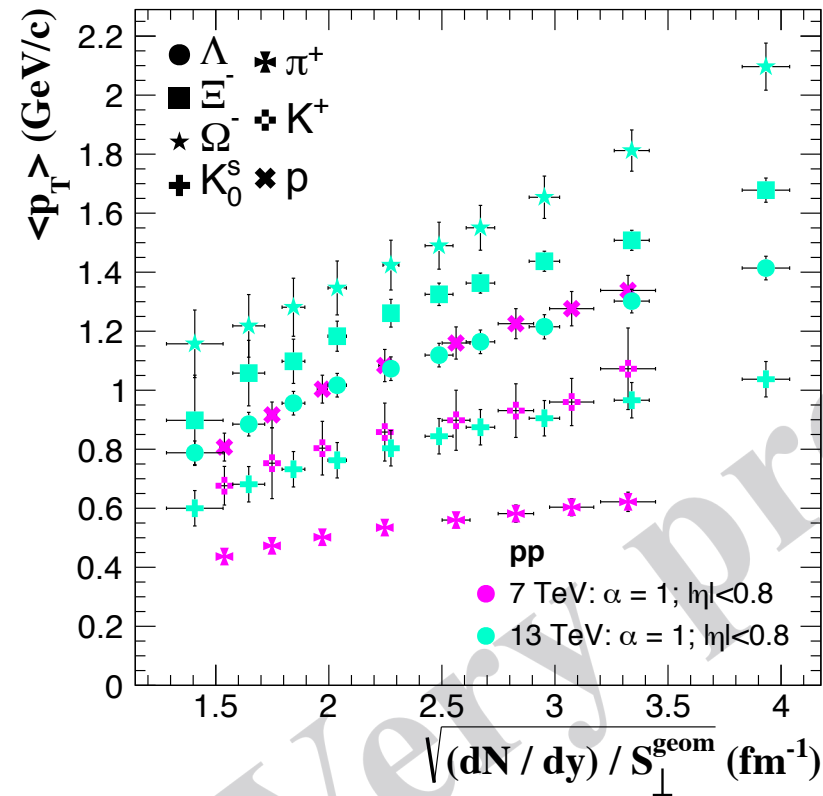
$\alpha=10$

$$f_{pp} = \begin{cases} -0.018 + 0.3976x + 0.095x^2 - 0.028x^3 & \text{if } x < 3.4 \\ 1.17 & \text{if } x \geq 3.4 \end{cases}$$

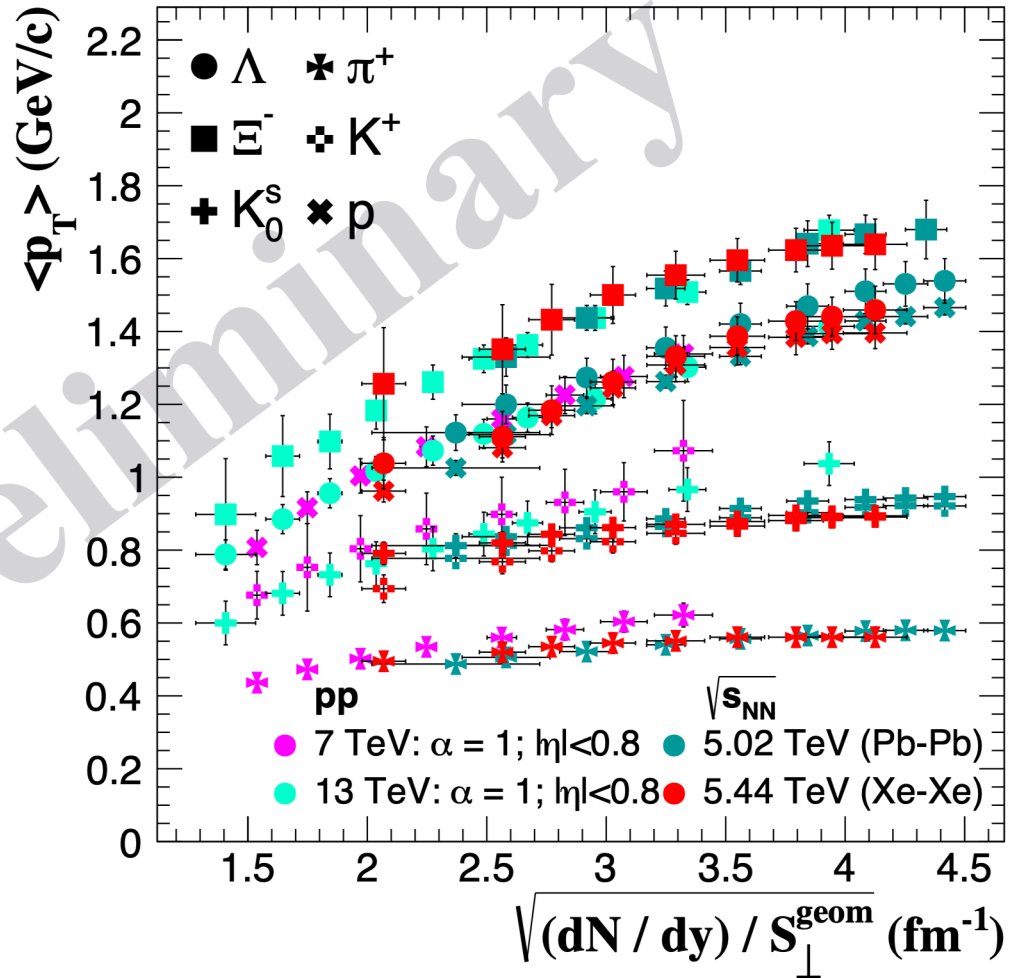
McLarren, M.Praszalowicz and B.Schenke, Nucl.Phys. A916(2013)210
 A.Bzdak, B.Schenke, P.Tribedy and R.Venugopalan, Phys.Rev. C87(2013)064906
 ALICE Collaboration, Nucl.Phys. A931(2014)c888

M. Petrovici, A. Lindner, A. Pop, M. Tartzila and I. Berceanu, Phys. Rev. C 98 (2018) 024904

p+p 13 TeV @ LHC



p+p 7 TeV vs. 13 TeV



p+p 7 TeV & 13 TeV vs. Pb+Pb & Xe+Xe @ LHC

Conclusions

- A very good scaling of different observables as a function of a scaling parameter suggested by the CGC model is found for Cu-Cu, Au-Au Xe-Xe and Pb-Pb for a wide range of energies: *from 7.7 GeV up to 5.44 TeV*
- The scaling is also evidenced for p+p and Pb+Pb at the measured LHC energies

Conclusions

- A very good scaling of different observables as a function of a scaling parameter suggested by the CGC model is found for Cu-Cu, Au-Au Xe-Xe and Pb-Pb for a wide range of energies: *from 7.7 GeV up to 5.44 TeV*
- The scaling is also evidenced for p+p and Pb+Pb at the measured LHC energies



The global trends at the LHC depend on the properties of the flux tubes of $\sim 1/[(dN/dy)/S_{\text{perp}}]^{1/2}$ size

Conclusions

- A very good scaling of different observables as a function of a scaling parameter suggested by the CGC model is found for Cu-Cu, Au-Au Xe-Xe and Pb-Pb for a wide range of energies: *from 7.7 GeV up to 5.44 TeV*
- The scaling is also evidenced for p+p and Pb+Pb at the measured LHC energies



The global trends at the LHC depend on the properties of the flux tubes of $\sim 1/[(dN/dy)/S_{\text{perp}}]^{1/2}$ size



System size is playing a minor role at LHC energies

Conclusions

- A very good scaling of different observables as a function of a scaling parameter suggested by the CGC model is found for Cu-Cu, Au-Au Xe-Xe and Pb-Pb for a wide range of energies: *from 7.7 GeV up to 5.44 TeV*
- The scaling is also evidenced for p+p and Pb+Pb at the measured LHC energies



The global trends at the LHC depend on the properties of the flux tubes of $\sim 1/[(dN/dy)/S_{\text{perp}}]^{1/2}$ size



System size is playing a minor role at LHC energies

Perspectives

- This study will be extended to pp collisions measured at LHC at 13 TeV - preliminary results already obtained
- The strange and multistrange hadrons seems to behave different in A-A relative to pp collisions - detailed analysis and interpretation is in progress

Thank you!