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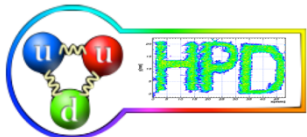
***Core-Corona
Geometrical scaling from RHIC to LHC energies
pp - (A-A) similarities***

Mihai Petrovici

Work done in collaboration with:

C.Andrei, I.Berceanu, A.Lindner, A.Pop, M.Tarzila, V.Topor Pop

HADRON PHYSICS DEPARTMENT



LIGHT UP Workshop, CERN, June 14-16, 2018

Outline

- **Physics motivation**
- **Core-corona interplay - impact on the experimental trends**
- **Geometrical scaling**

Au-Au @ RHIC, Pb-Pb @ LHC

- $\langle p_T \rangle$ vs. $[(dN/dy)/S_{perp}]^{1/2}$

- *The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{perp}]^{1/2}$*

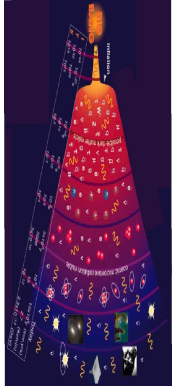
- $\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{perp}]^{1/2}$

Au-Au & Cu-Cu @ RHIC, Pb-Pb & Xe-Xe @ LHC

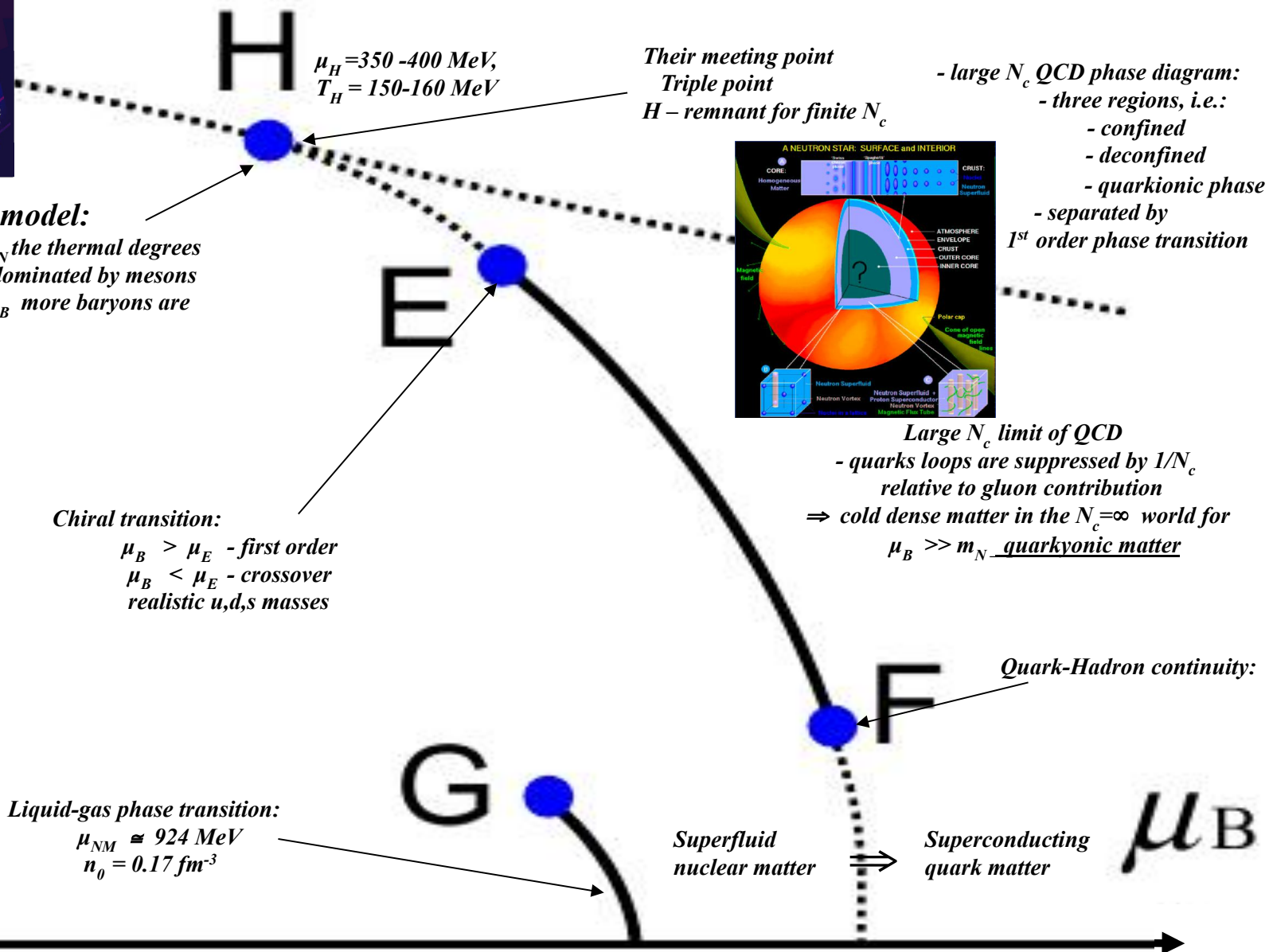
p+p vs. Pb-Pb @ LHC

- **Outlook**

Physics motivation



T



Statistical model:

- for $\mu_B \ll m_N$ the thermal degrees of freedom dominated by mesons
- for higher μ_B more baryons are excited

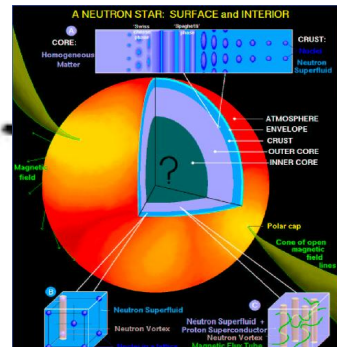
Chiral transition:
 $\mu_B > \mu_E$ - first order
 $\mu_B < \mu_E$ - crossover
 realistic u, d, s masses

Liquid-gas phase transition:
 $\mu_{NM} \approx 924 \text{ MeV}$
 $n_0 = 0.17 \text{ fm}^{-3}$

H
 $\mu_H = 350 - 400 \text{ MeV}$
 $T_H = 150 - 160 \text{ MeV}$

Their meeting point
 Triple point
 H – remnant for finite N_c

- large N_c QCD phase diagram:
- three regions, i.e.:
- confined
- deconfined
- quarkionic phase
- separated by 1st order phase transition



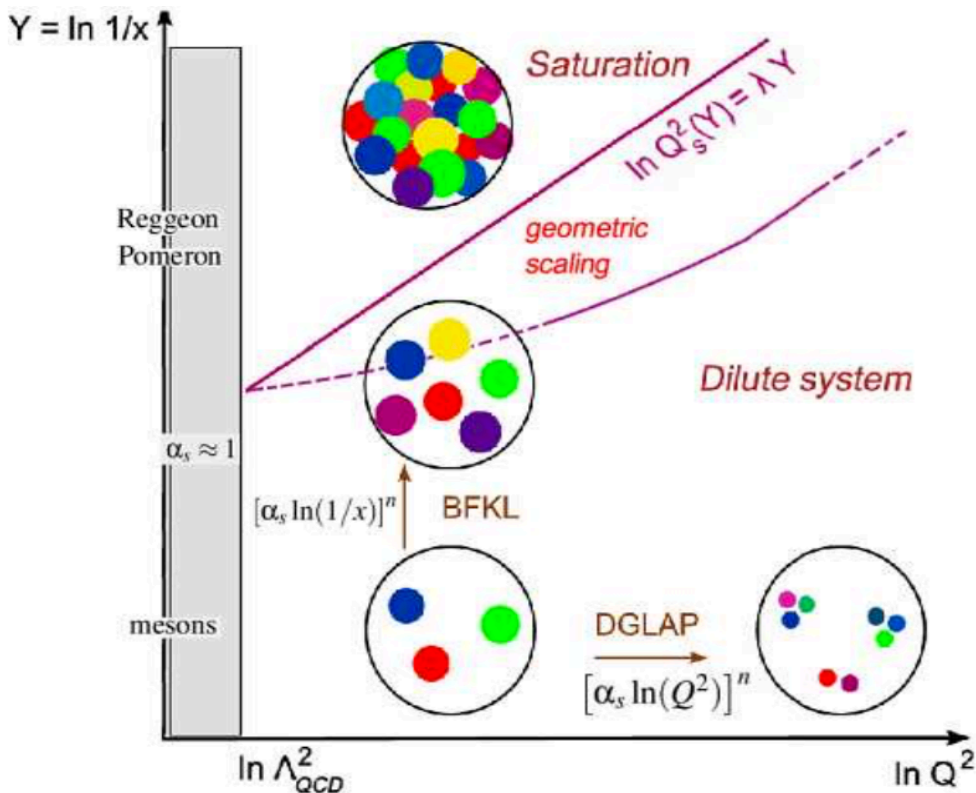
Large N_c limit of QCD
 - quarks loops are suppressed by $1/N_c$ relative to gluon contribution
 \Rightarrow cold dense matter in the $N_c = \infty$ world for $\mu_B \gg m_N$ quarkyonic matter

Quark-Hadron continuity:

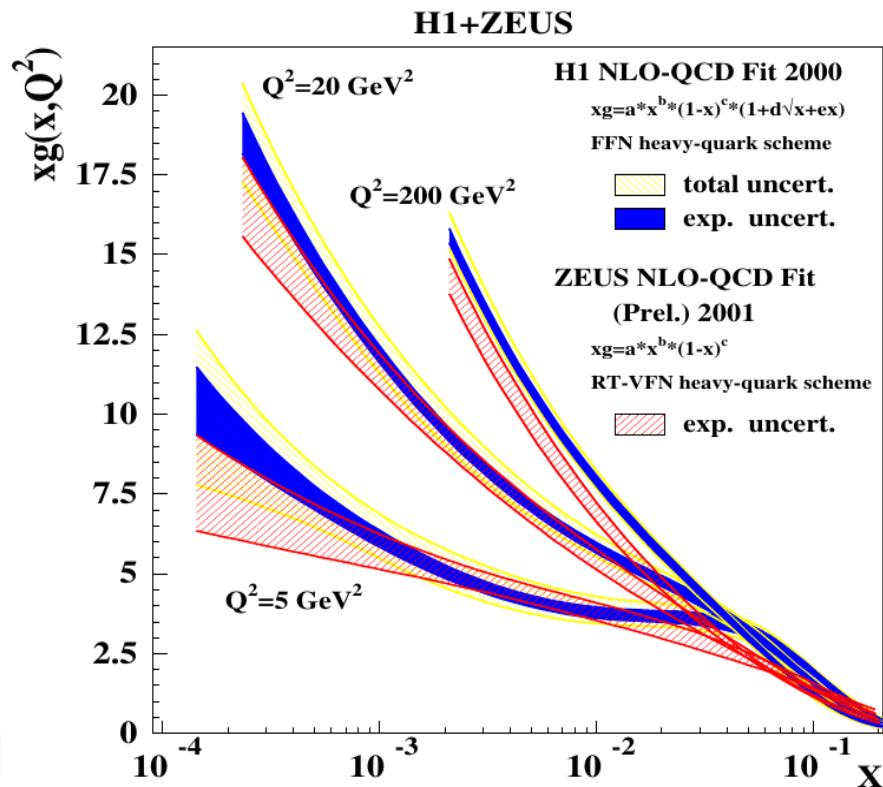
Superfluid nuclear matter \Rightarrow Superconducting quark matter

μ_B

Physics motivation



D. d'Enterria, Eur.Phys.J. A31(2007)816

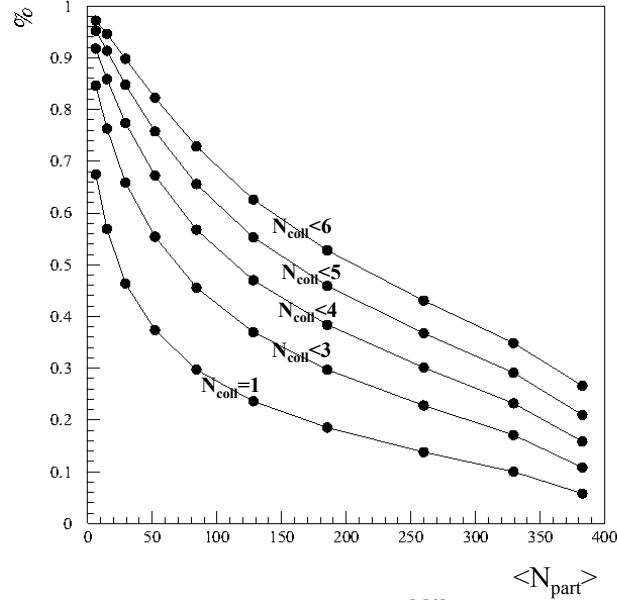
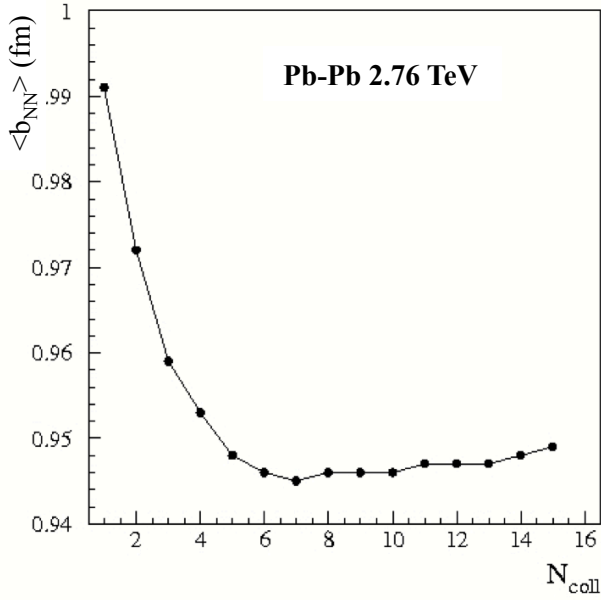


M.Dittmar et al., Proceedings HERA-LHC Workshop
arXiv:[hep-ph]0511119

System	Au-Au	Pb-Pb	Pb-Pb	pp
$\sqrt{s}(\text{GeV})$	200	2700	5020	7000
$\frac{dN_g^{in}}{dyd^2b}(\text{fm}^{-2})$	≈ 4.7	≈ 11.8	≈ 15.9	≈ 18.7
f_{in}^g	≈ 0.9	≈ 2.3	≈ 3.1	≈ 3.6

Following A.H. Mueller
approximations NP A715(2003)20

Core-Corona effect Glauber MC

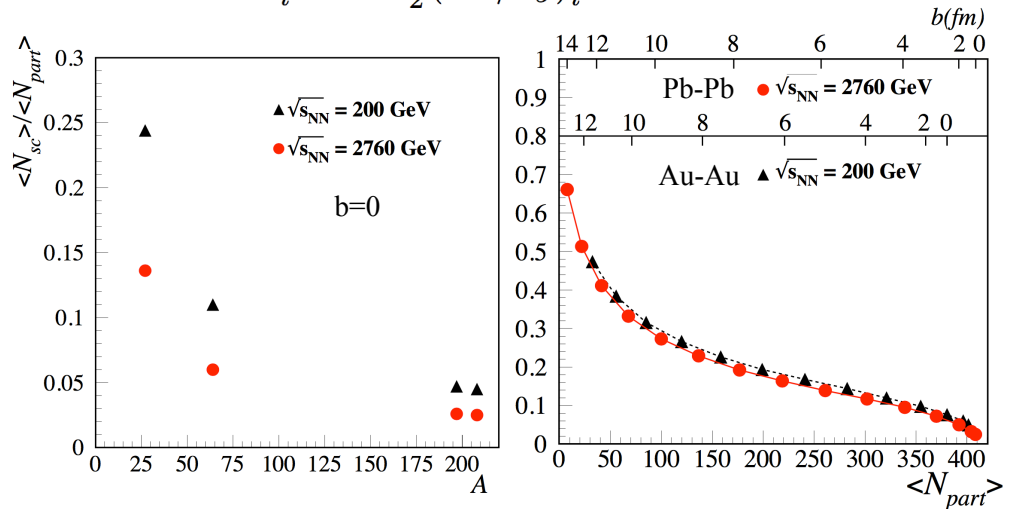


How should we define “Corona” ?

$$\left(\frac{dN}{dy}\right)_i^{cen} = N_{part} [(1 - f_{core}) M_i^{ppMB} + f_{core} M_i^{core}]$$

$$M_i^{ppMB} = \frac{1}{2} (dN/dy)_i^{ppMB}$$

F.Becattini, J.Manninen, J.Phys.G 35 (2008) 104013, Phys.Lett.B 673(2009)19
 J.Aichelin, K.Werner, Phys.Rev.C 79(2009)064907,
 J.Phys.G 37(2010)094006,
 Phys.Rev.C 82(2010)034906
 P.Bozek, Phys.Rev.C 79 (2009) 054901
 C.Schreiber, K.Werner, J.Aichelin, Phys.Atom.Nucl. 75(2012)640
 M.Gemard and J.Aichelin, Astron.Nachr. 335(2014)660

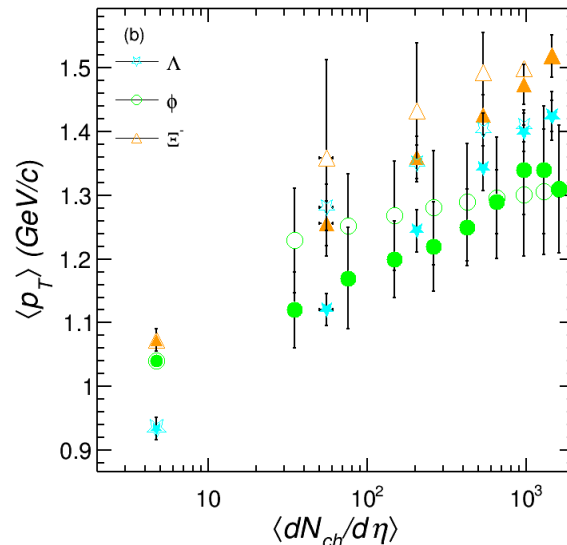
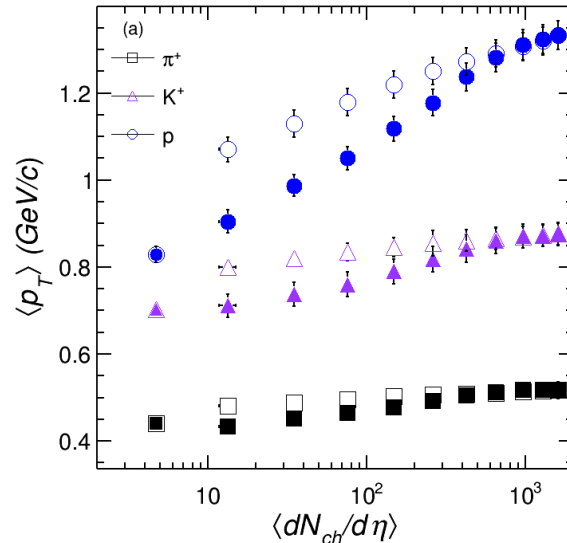
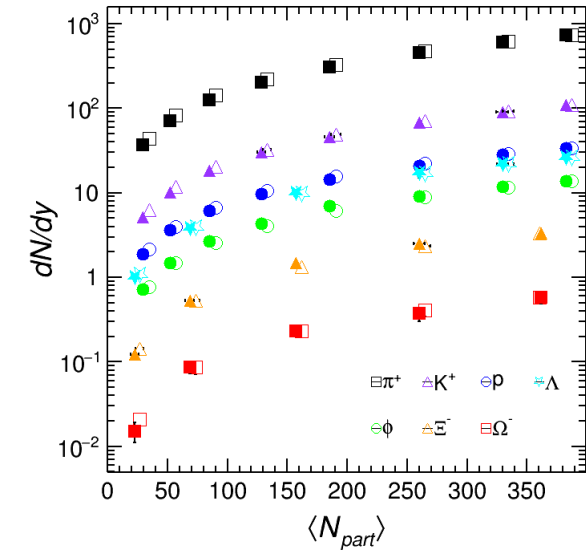


Core-Corona effect

$$\left(\frac{dN}{dy}\right)_i^{cen} = N_{part}[(1 - f_{core})M_i^{ppMB} + f_{core}M_i^{core}] \quad (1)$$

$$M_i^{ppMB} = \frac{1}{2}(dN/dy)_i^{ppMB} \quad (2)$$

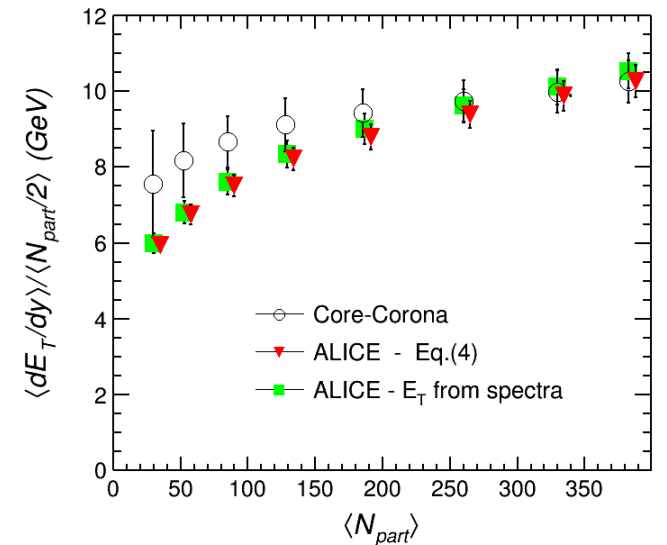
$$\langle p_T \rangle_i^{cen} = \frac{f_{core}\langle p_T \rangle_i^{core} M_i^{core} + (1 - f_{core})\langle p_T \rangle_i^{ppMB} M_i^{ppMB}}{f_{core}M_i^{core} + (1 - f_{core})M_i^{ppMB}} \quad (3)$$



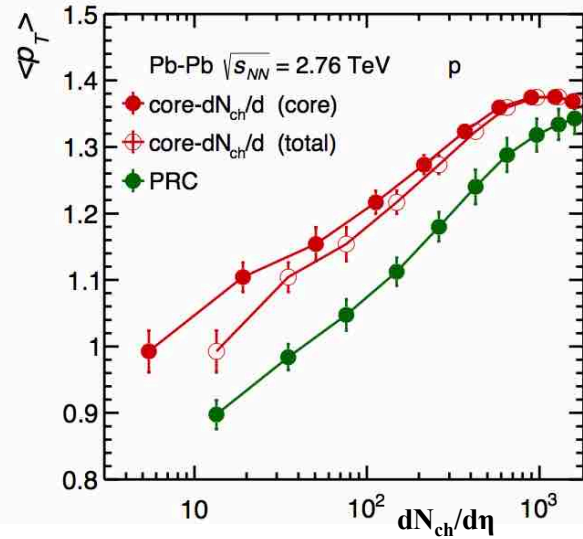
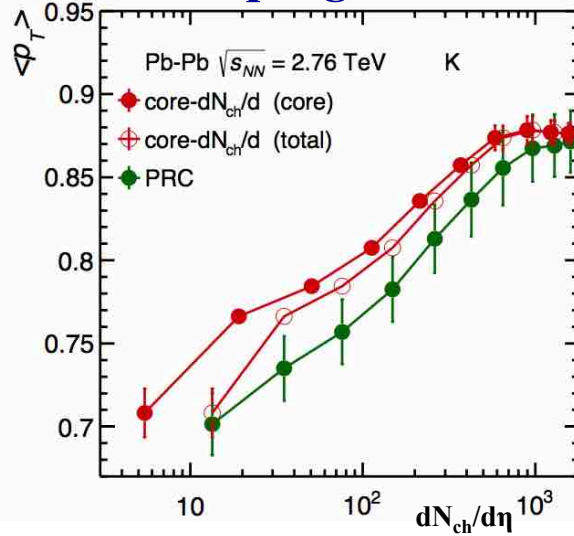
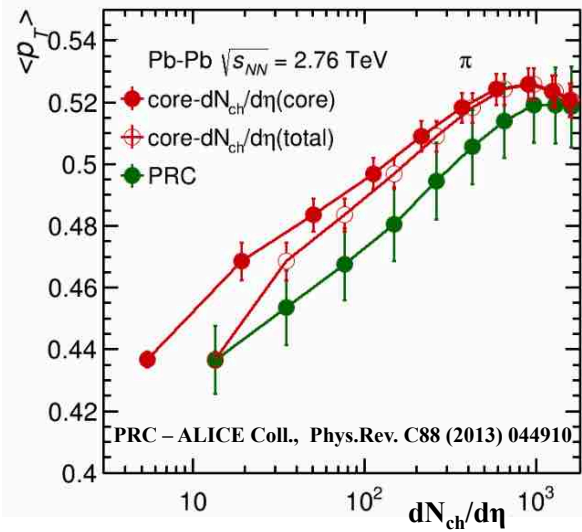
Full symbols - experiment
Open symbols - core-corona interplay

$$\frac{dE_T}{dy} \approx 3 \left(\frac{dE_T}{dy}\right)_{\pi^+} + 4 \left(\frac{dE_T}{dy}\right)_{K^+, p, \Xi^-} + 2 \left(\frac{dE_T}{dy}\right)_{\Lambda, \Omega^-} \quad (4)$$

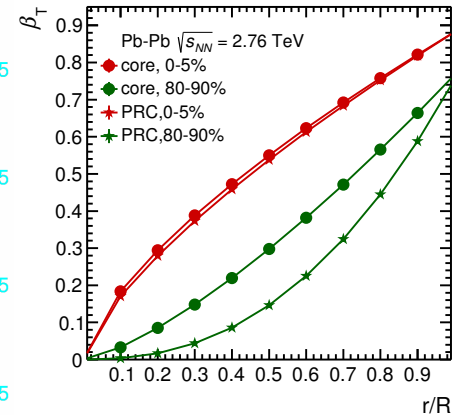
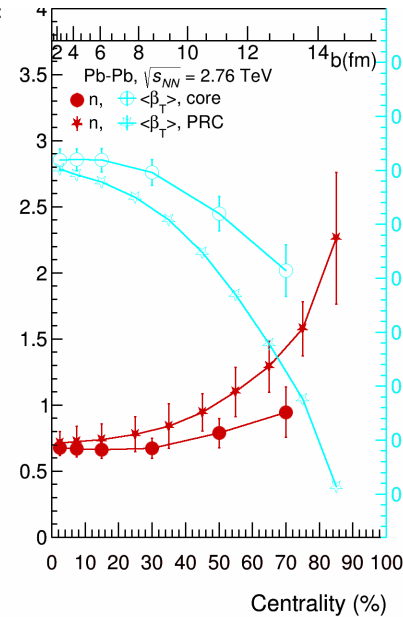
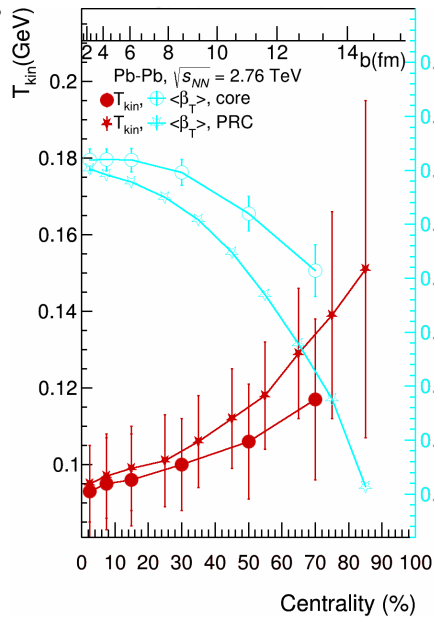
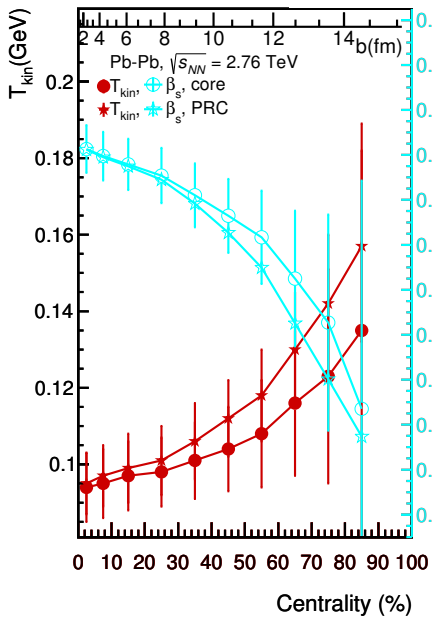
$$\frac{dE_T}{dy} = \langle m_T \rangle \frac{dN}{dy} \quad \langle m_T \rangle = \sqrt{\langle p_T \rangle^2 + m^2}$$



Core properties In progress



BGBW $\frac{1}{2\pi p_T} \frac{d^2N}{dy dp_T} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}} \right)$ $\rho = \tanh^{-1} \beta_T = \tanh^{-1} \left[\left(\frac{r}{R} \right)^n \beta_s \right]$



Based on ALICE data, see references in:
M. Petrovici, I. Berceanu, A. Pop, M. Târzila, and C. Andrei, Phys.Rev. C96(2017)014908

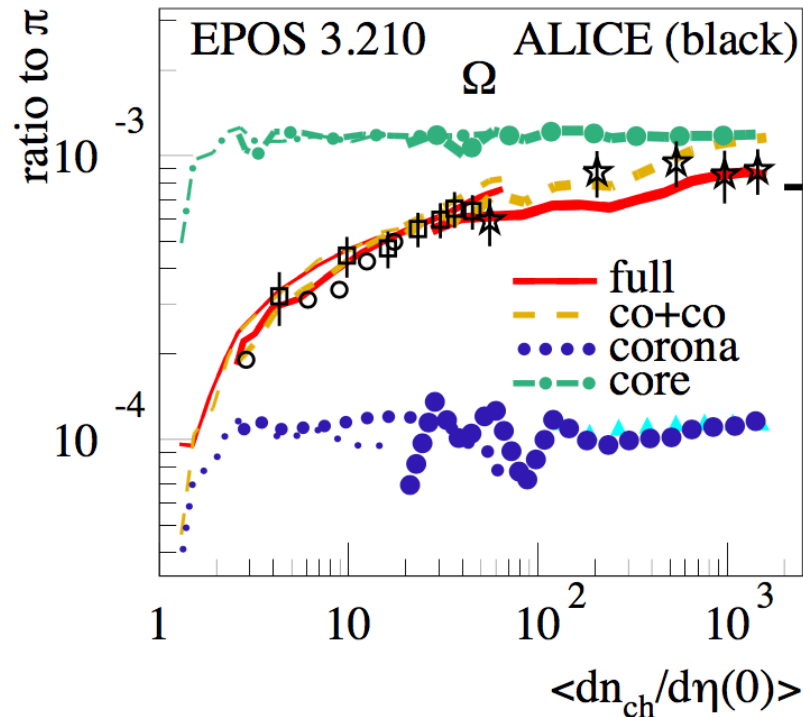
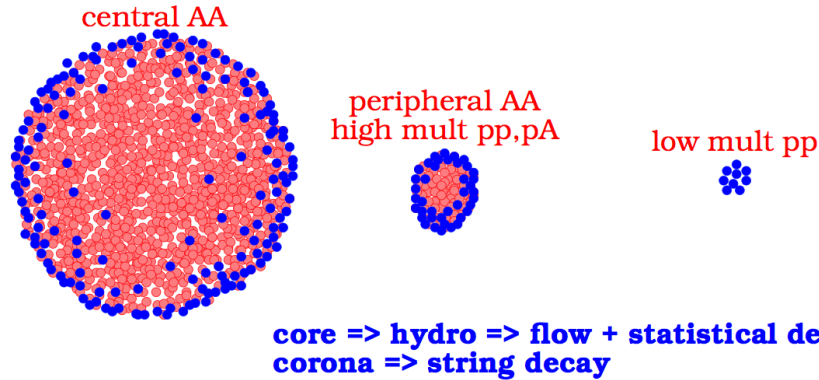
Core-corona in pp

Core-corona picture in EPOS

Phys.Rev.Lett. 98 (2007) 152301, Phys.Rev. C89 (2014) 6, 064903

K. Werner, SQM 2017, July 10-15 2017, Utrecht

Gribov-Regge approach => (Many) kinky strings
=> core/corona separation (based on string segments)



Geometrical scaling

*Local parton-hadron duality picture
and dimensionality argument*



$$\langle p_T \rangle / \sqrt{\frac{dN}{dy} / S_{\perp}} \sim \frac{1}{n\sqrt{n}}$$

n - no. of charged part. from a gluon fragmentation



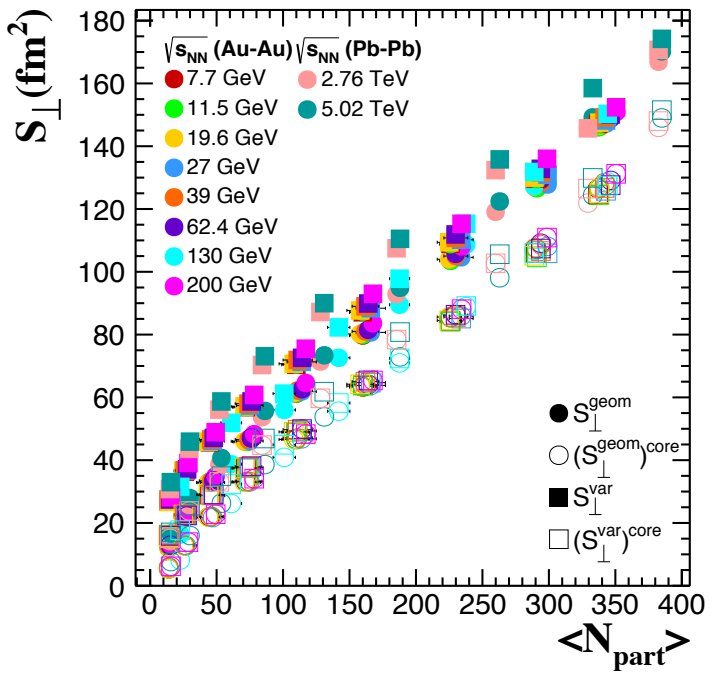
$$\langle p_T \rangle / \sqrt{\frac{dN}{dy} / S_{\perp}}$$

decreases as a function of:
- collision energy
- centrality

- Y.L.Dokshitzer, V.A.Khoze and S.Troian, J.Phys.G 17 (1991) 1585
- T. Lappi, Eur.Phys.J. C71 (2011) 1699
- E. Levin and A.H. Rezaian, Phys.Rev.D 83 (2011) 114001

S_{perp} & dN/dy estimates

Glauber Monte Carlo approach



$$S_{\perp}^{\text{var}} \sim \pi \sqrt{\sigma_x^2 \sigma_y^2 - \sigma_{xy}^2}$$

BES

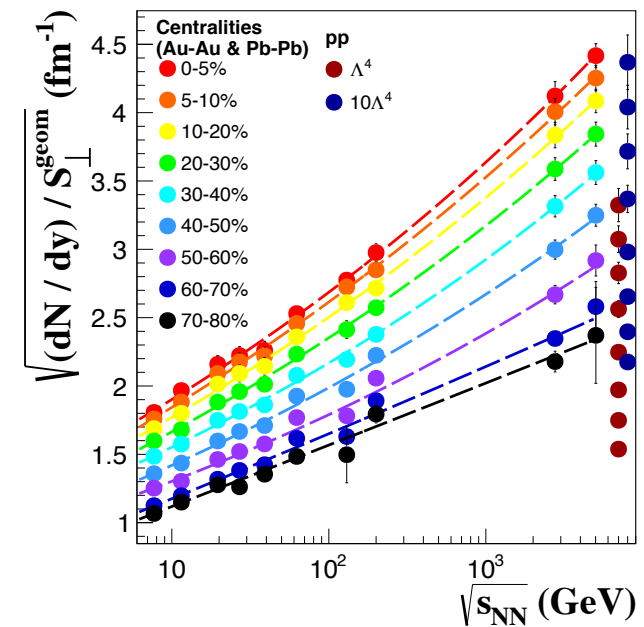
RHIC
62.4; 130; 200 GeV

LHC

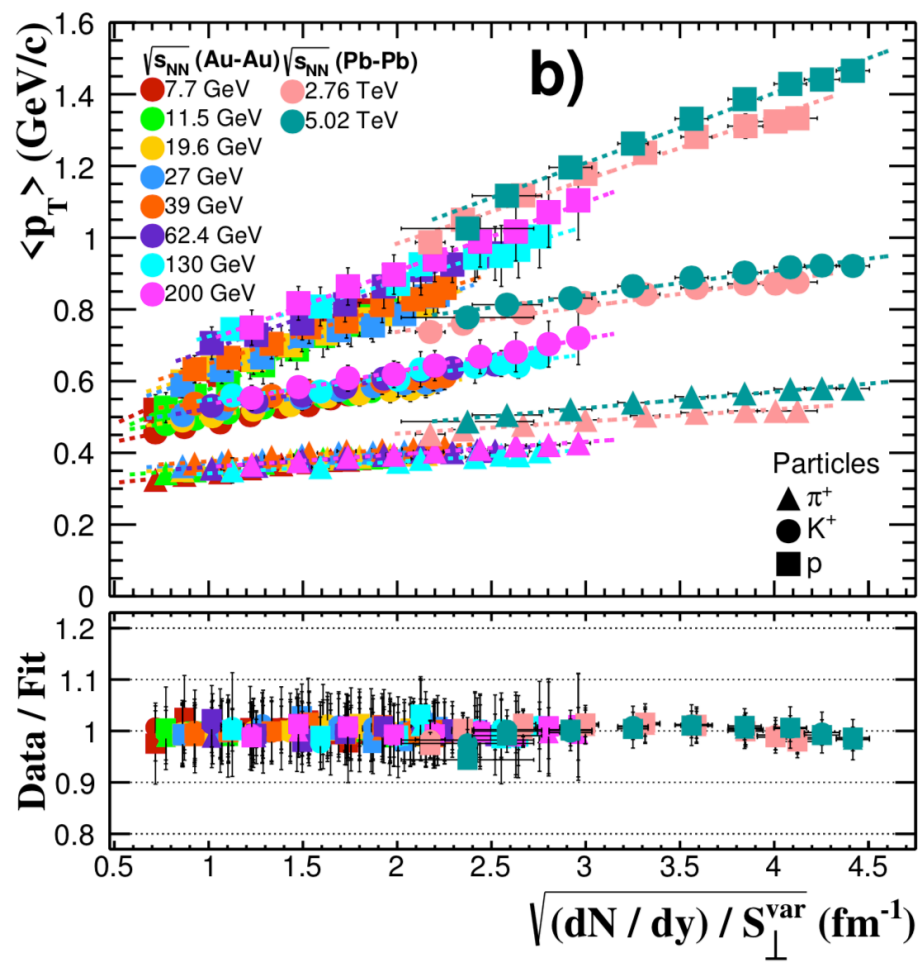
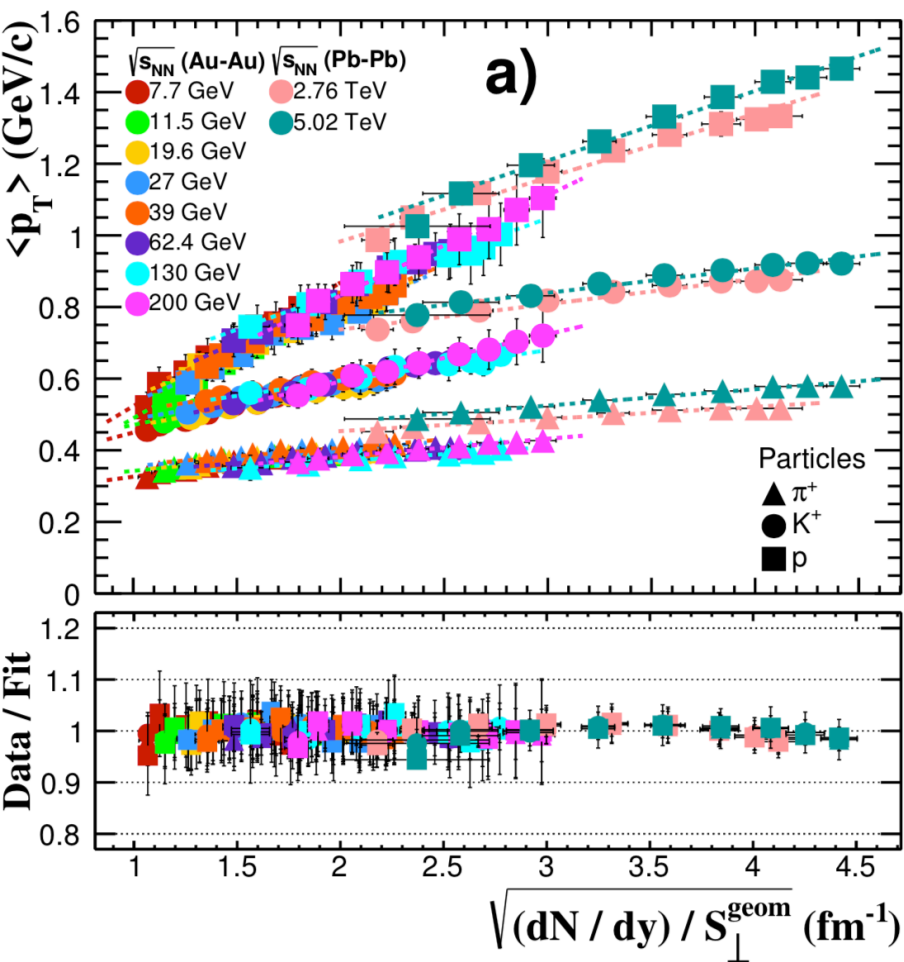
$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN(\pi^+ + \pi^-)}{dy} + 2 \frac{dN(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)}{dy} + \frac{dN(\Lambda + \bar{\Lambda})}{dy}$$

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN(\pi^+ + \pi^-)}{dy} + 2 \frac{dN(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)}{dy} + \frac{dN(\Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}{dy}$$

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN(\pi^+ + \pi^-)}{dy} + 2 \frac{dN(p + \bar{p}, \Xi^- + \bar{\Xi}^+)}{dy} + \frac{dN(K^+ + K^-, K_S^0, \Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}{dy}$$

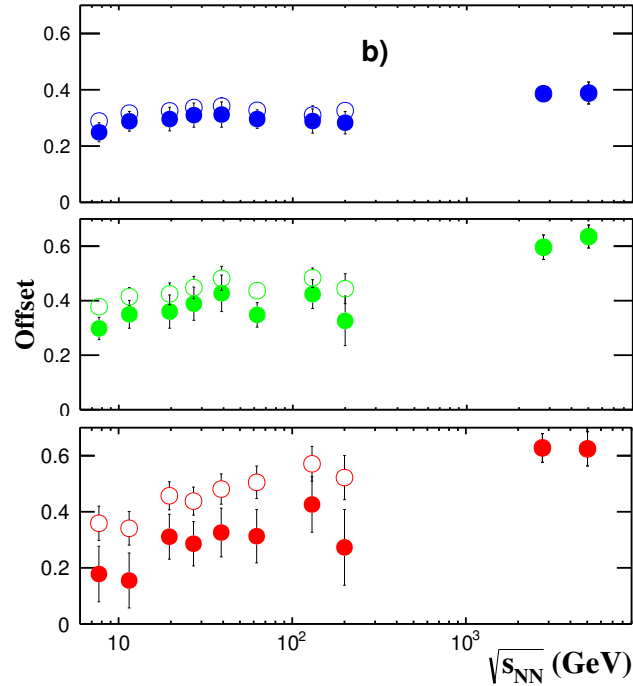
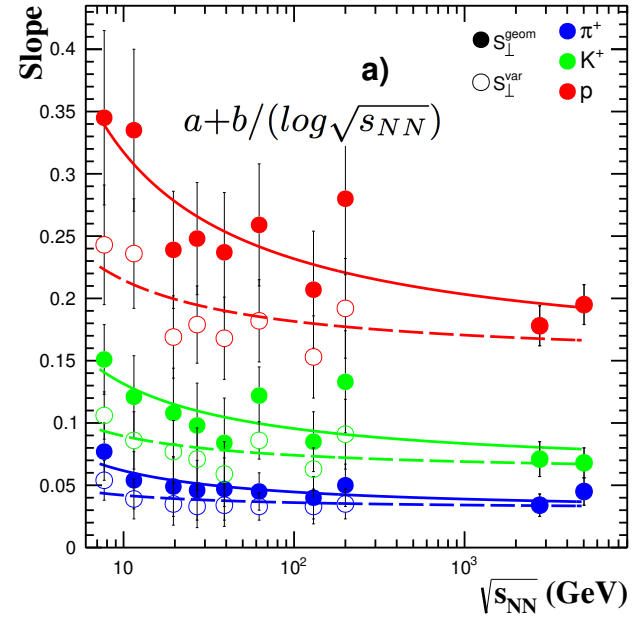


$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



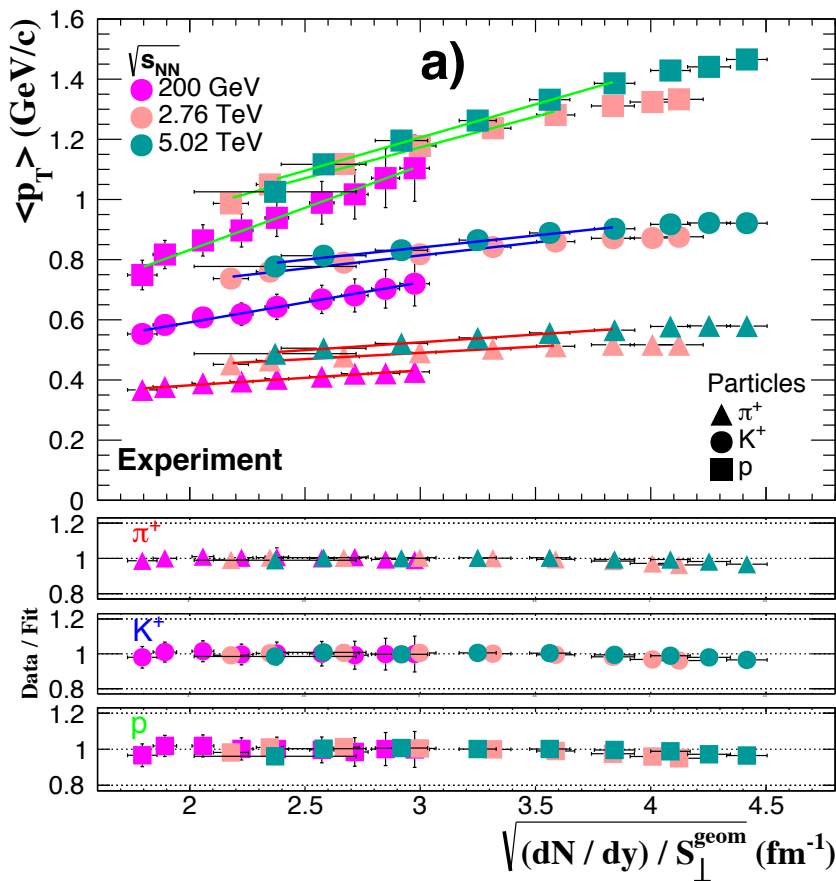
STAR Collaboration, Phys.Rev. C96(2017)044904
 STAR Collaboration, Phys.Rev. C79(2009)034909
 ALICE Collaboration, Phys.Rev. C{88}{2013}{044910}
 ALICE Collaboration, Phys.Rev.Lett. 116(2016)222302
 ALICE Collaboration, Eur.Phys.J. C75(2015)226
 A.K.Dash, ALICE Collaboration, 9th Int. Workshop on MPI at LHC, Dec. 11-15, 2017

$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$

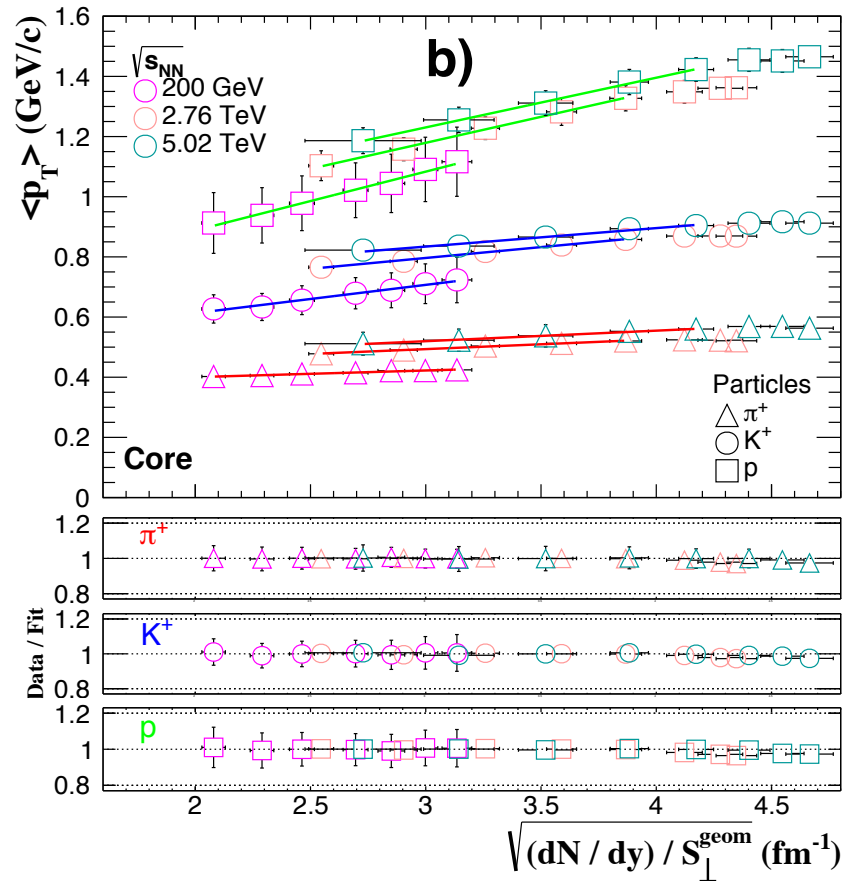


$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$

Core-Corona effect



$\sqrt{s_{NN}}$ (GeV)	Slope			Offset		
	π^+	K^+	p	π^+	K^+	p
200	0.05 ± 0.02	0.13 ± 0.04	0.28 ± 0.06	0.28 ± 0.04	0.33 ± 0.09	0.27 ± 0.13
2760	0.04 ± 0.01	0.09 ± 0.02	0.20 ± 0.03	0.37 ± 0.04	0.56 ± 0.07	0.56 ± 0.08
5020	0.05 ± 0.02	0.08 ± 0.02	0.22 ± 0.03	0.37 ± 0.06	0.60 ± 0.07	0.54 ± 0.10



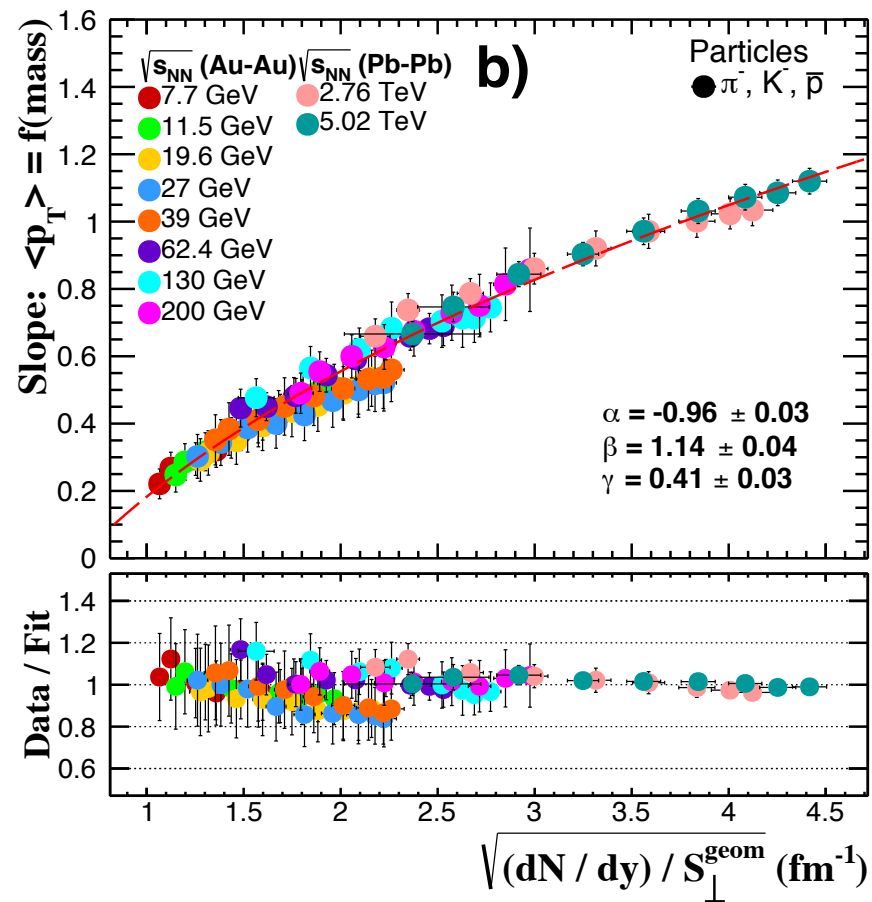
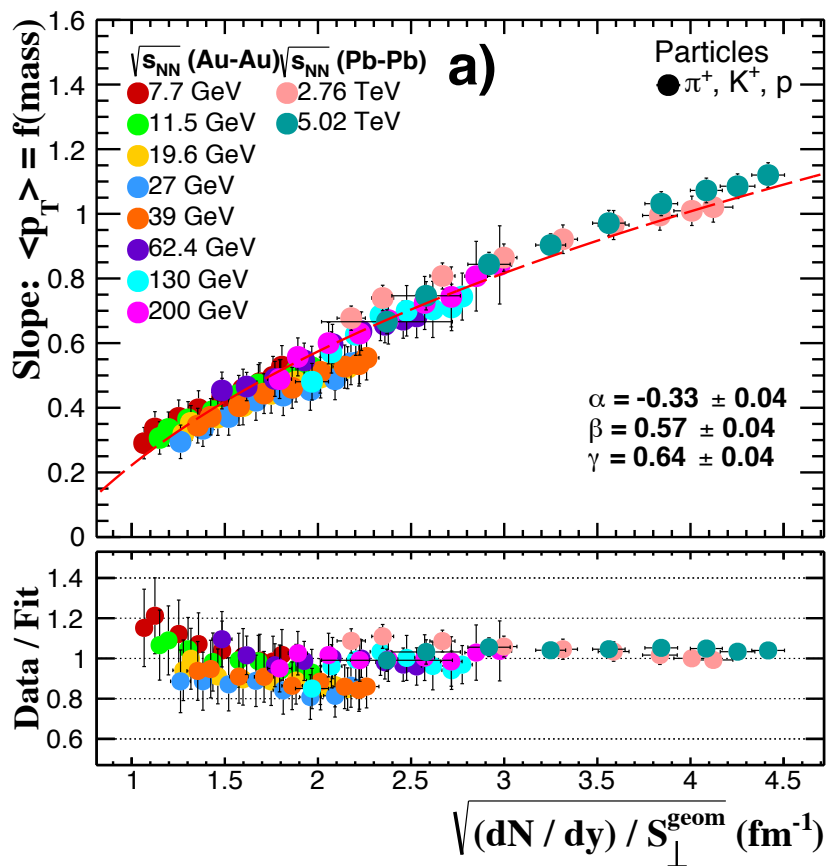
$\sqrt{s_{NN}}$ (GeV)	Slope			Offset		
	π^+	K^+	p	π^+	K^+	p
200	0.02 ± 0.03	0.09 ± 0.06	0.20 ± 0.11	0.36 ± 0.07	0.43 ± 0.15	0.50 ± 0.29
2760	0.03 ± 0.02	0.07 ± 0.03	0.17 ± 0.04	0.40 ± 0.06	0.58 ± 0.10	0.66 ± 0.14
5020	0.03 ± 0.03	0.06 ± 0.02	0.17 ± 0.04	0.41 ± 0.11	0.65 ± 0.08	0.73 ± 0.16

$$\langle p_T \rangle_i^{\text{cen}} = \frac{f_{\text{core}} \langle p_T \rangle_i^{\text{core}} M_i^{\text{core}} + (1 - f_{\text{core}}) \langle p_T \rangle_i^{\text{ppMB}} M_i^{\text{ppMB}}}{f_{\text{core}} M_i^{\text{core}} + (1 - f_{\text{core}}) M_i^{\text{ppMB}}}$$

$$\left(\frac{dN}{dy} \right)_i^{\text{cen}} = \langle N_{\text{part}} \rangle [(1 - f_{\text{core}}) M_i^{\text{ppMB}} + f_{\text{core}} M_i^{\text{core}}]$$

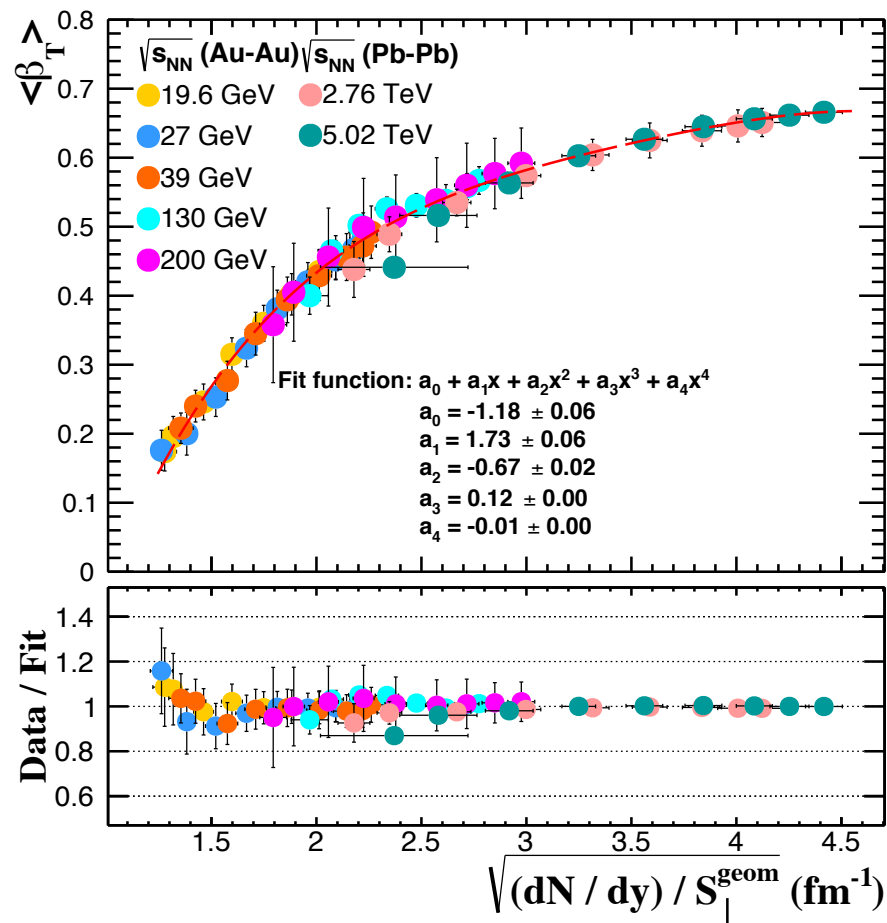
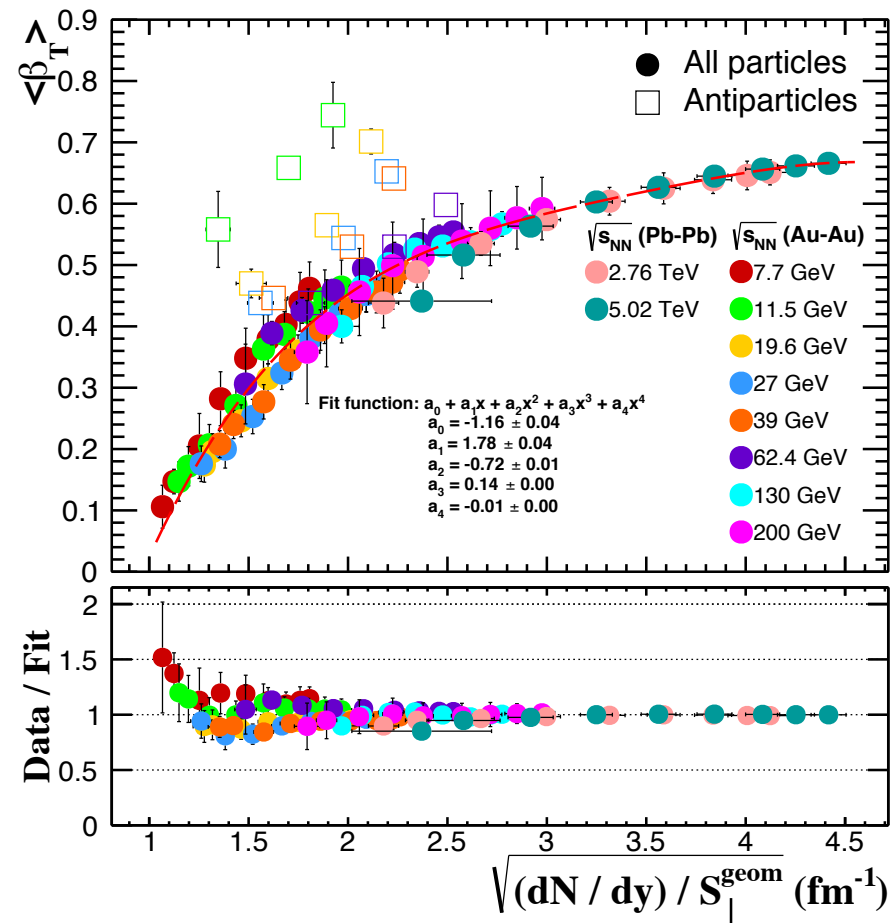
$$M_i^{\text{ppMB}} = \frac{1}{2} (dN/dy)_i^{\text{ppMB}}$$

The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



$$\text{Slope}_{\langle p_T \rangle = f(\text{mass})} = \alpha + \beta \left(\sqrt{\frac{dN}{dy} / S_{\perp}^{\text{geom}}} \right)^{\gamma}$$

$\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$

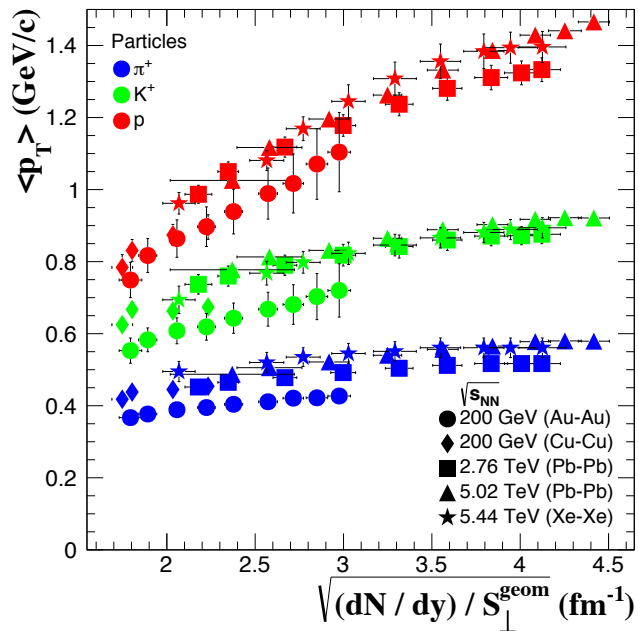


Boltzmann-Gibbs
Blast Wave

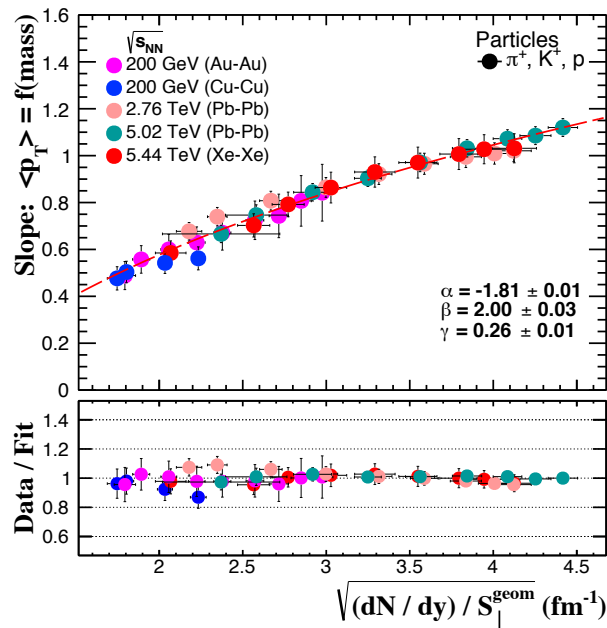
$$\left\{ \begin{array}{l} \frac{1}{2\pi p_T} \frac{d^2 N}{dy dp_T} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}} \right) \\ \rho = \tanh^{-1} \beta_T = \tanh^{-1} \left[\left(\frac{r}{R} \right)^n \beta_s \right] \end{array} \right.$$

Cu-Cu; Au-Au @ RHIC vs. Xe-Xe and Pb+Pb @ LHC

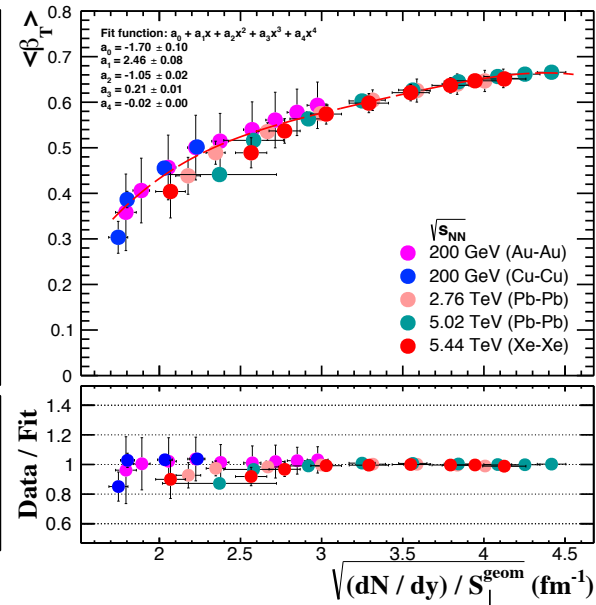
$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



$\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



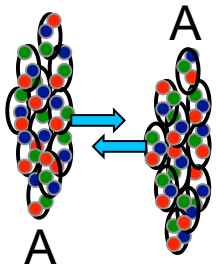
BRAHMS Collaboration, arXic:[nucl.ex]1602.01183

F.Bellini, ALICE Collaboration, QM2018

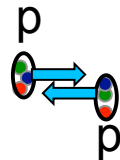
STAR Collaboration, Phys.Rev. C79(2009)034909

ALICE Collaboration, Phys.Rev. C}{88}{2013}{044910}

p+p vs. Pb+Pb @ LHC

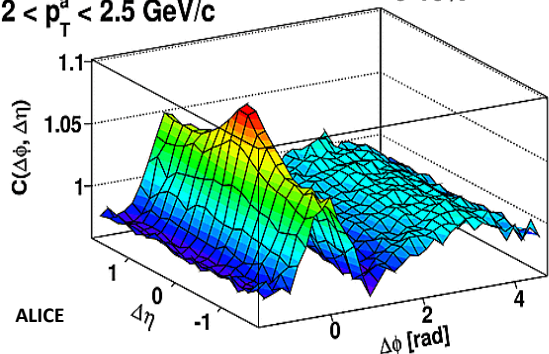


Long range near side correlations



$3 < p_T^t < 4 \text{ GeV}/c$
 $2 < p_T^a < 2.5 \text{ GeV}/c$

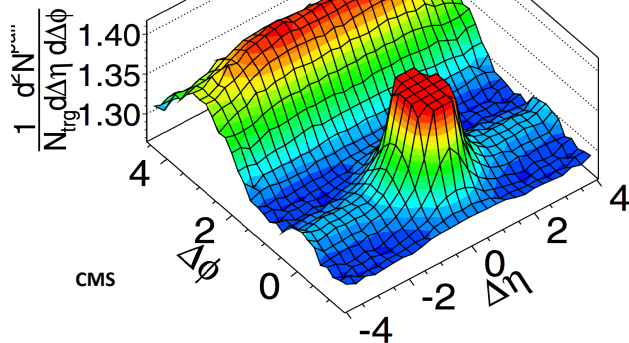
Pb-Pb 2.76 TeV
 0-10%



Phys.Lett. B708(2014)249

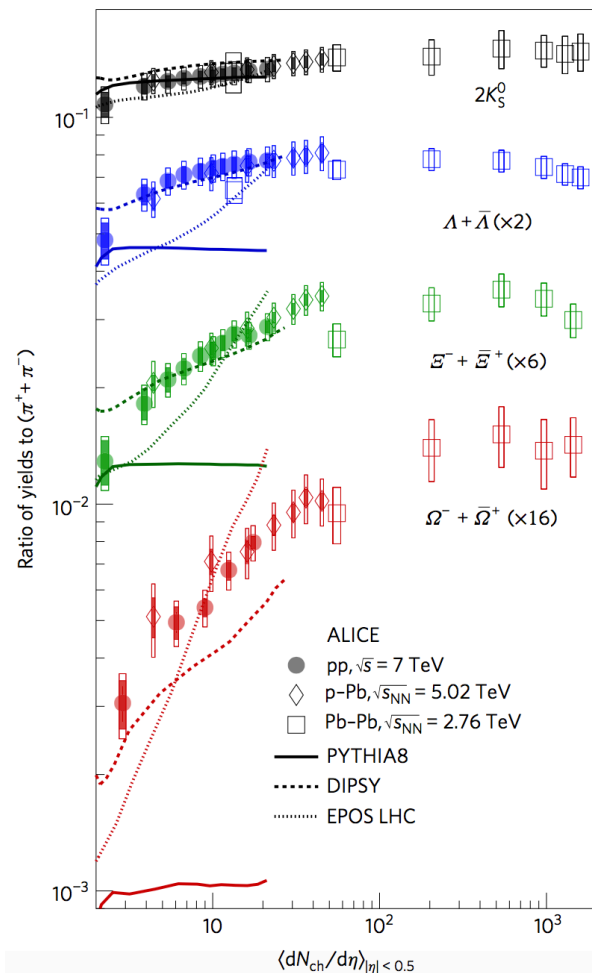
pp $\sqrt{s} = 7 \text{ TeV}$, $N \geq 110$ $|\eta| < 2.4$

$2 < p_T^{\text{trig}} < 3 \text{ GeV}/c$
 $1 < p_T^{\text{assoc}} < 2 \text{ GeV}/c$

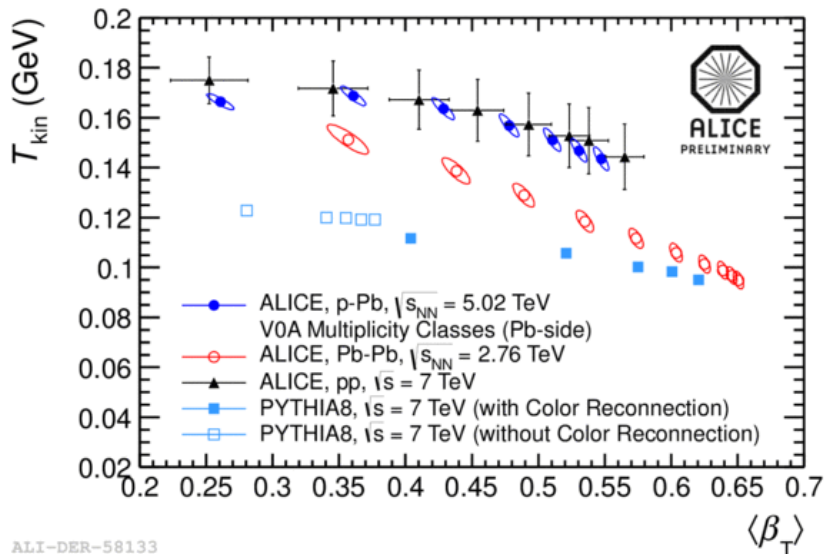


J. Phys. G Nucl. Part. Phys. 38(2011)124051

Strangeness relative yields



Transverse flow – BGBW fits



ALICE Collaboration, Nucl. Phys. A 931 (2014) c888

ALICE Collaboration, Nat. Phys. 13(2017)535

p+p vs. Pb+Pb @ LHC

$$S_{\perp}^{pp} = \pi R_{pp}^2 \quad R_{pp} = l_{fm} \mathbf{f}_{pp} - \text{maximal radius for which the energy density of the Yang-Mill fields is larger than } \varepsilon = \alpha \Lambda_{QCD}^4 \quad (\alpha \in [1, 10])$$

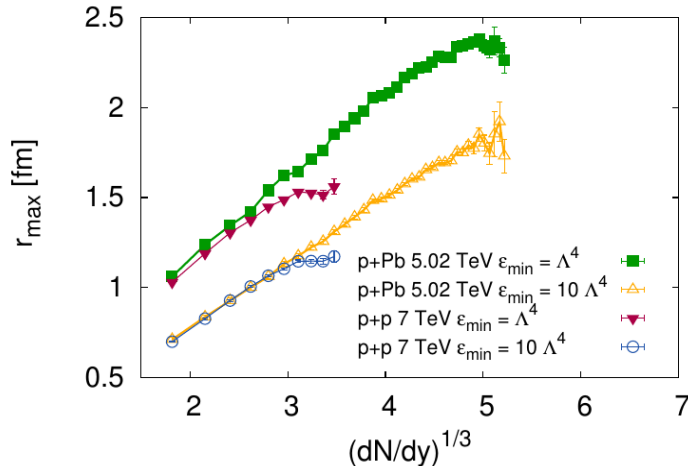
$$\alpha=1 \quad f_{pp} = \begin{cases} 0.387 + 0.0335x + 0.274x^2 - 0.0542x^3 & \text{if } x < 3.4 \\ 1.538 & \text{if } x \geq 3.4 \end{cases}$$

$$x = (dN_g/dy)^{1/3}$$

$$dN_g/dy \approx dN/dy$$

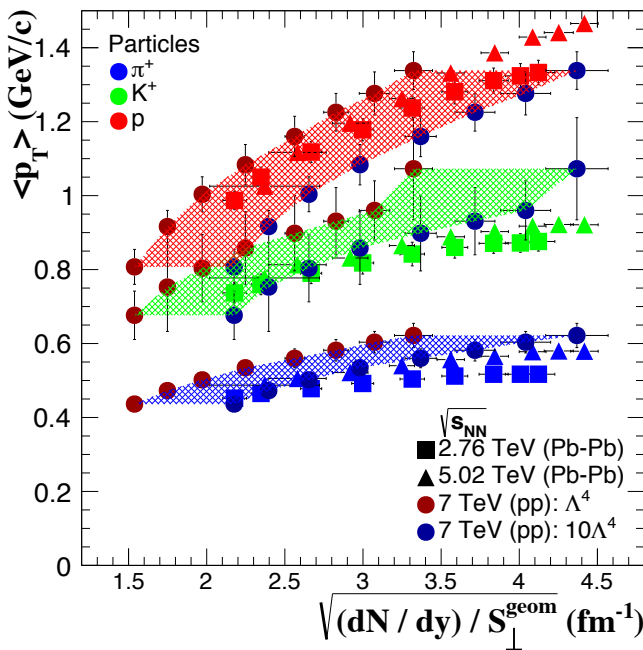
McLarren, M.Praszalowicz and B.Schenke, Nucl.Phys. A916(2013)210

$$\alpha=10 \quad f_{pp} = \begin{cases} -0.018 + 0.3976x + 0.095x^2 - 0.028x^3 & \text{if } x < 3.4 \\ 1.17 & \text{if } x \geq 3.4 \end{cases}$$

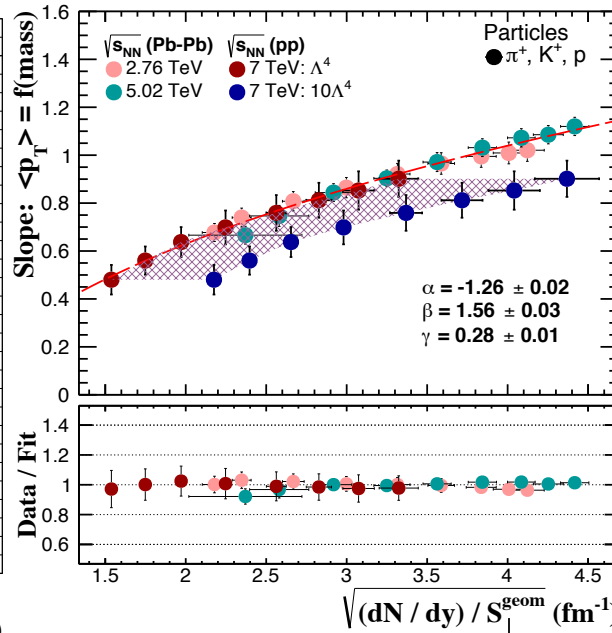


A.Bzdak, B.Schenke, P.Tribedy and R.Venugopalan, Phys.Rev. C87(2013)064906

$\langle p_T \rangle$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$

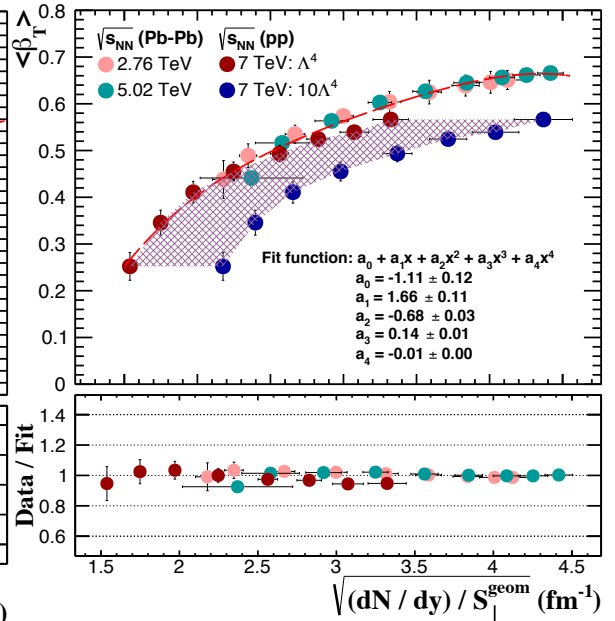


The slope of $\langle p_T \rangle = f(\text{mass})$ vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



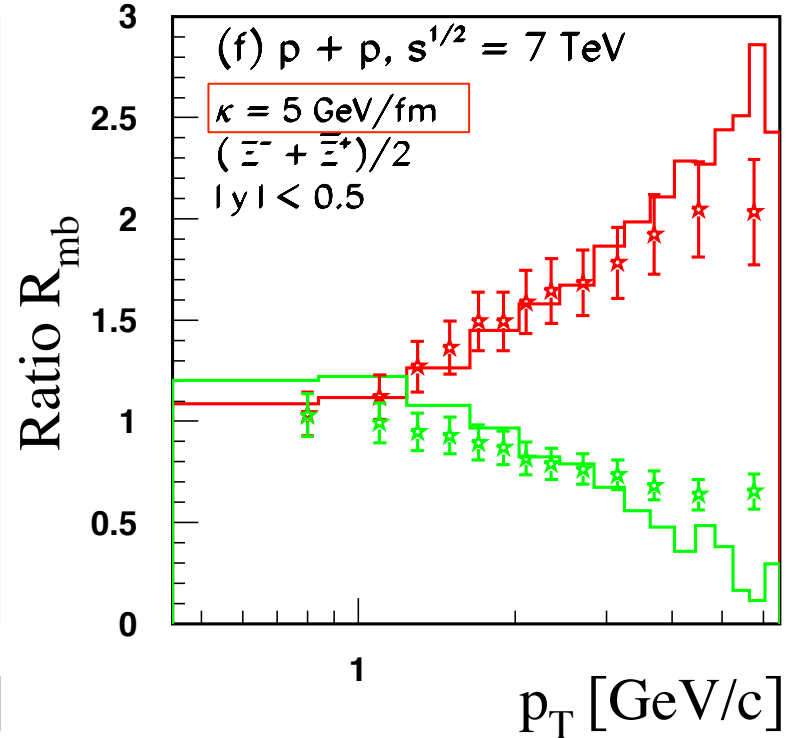
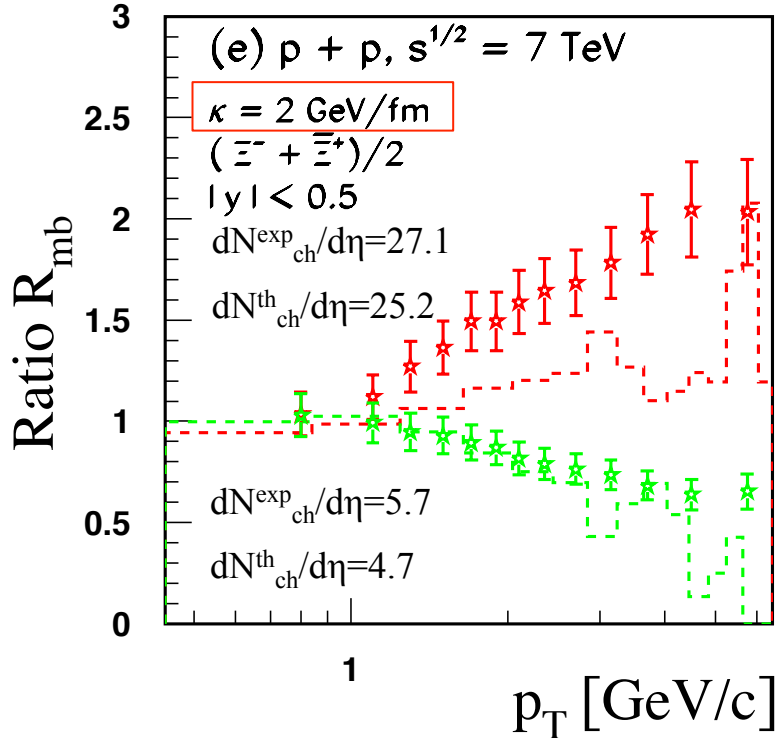
$$\text{Slope}_{\langle p_T \rangle = f(\text{mass})} = \alpha + \beta \left(\sqrt{\frac{dN}{dy} / S_{\perp}^{\text{geom}}} \right)^{\gamma}$$

$\langle \beta_T \rangle$ from BGBW fits vs. $[(dN/dy)/S_{\text{perp}}]^{1/2}$



pp - Pb+Pb similarities @ LHC within HIJING/BB v2.0 model

$$R_{mb} (cen) = \left(\frac{d^2 N}{dy dp_T} \right)^{cen} / \left(\frac{d^2 N}{dy dp_T} \right)^{ppMB} / \left(\frac{\langle dN_{ch} \rangle}{d\eta} \right)_i$$



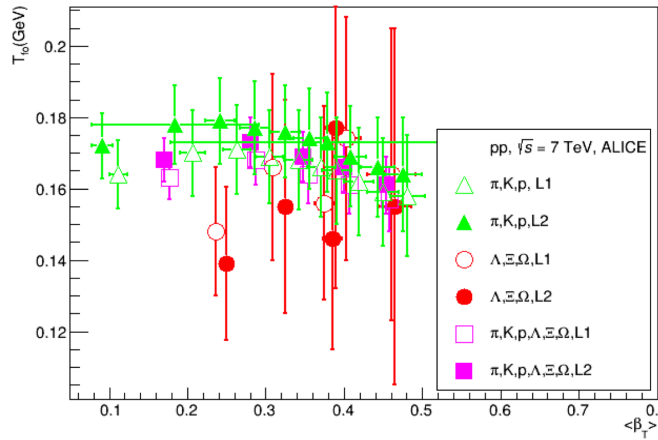
V.Topor Pop and M.Petrovici, arXiv:[hep-ph]1806.00359

ALICE Collaboration, Phys.Lett. 712B(2012)309

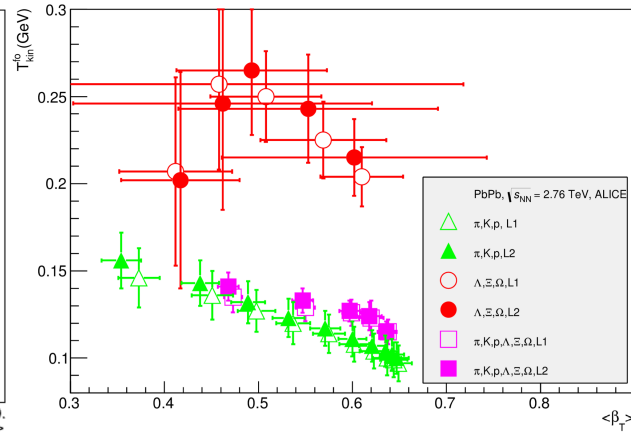
R. Derradi de Souza, ALICE Collaboration, J. Phys. Conf. Ser. 779, no. 1(2017)012071

There are still some differences between pp & $A-A$ BGBW - fits

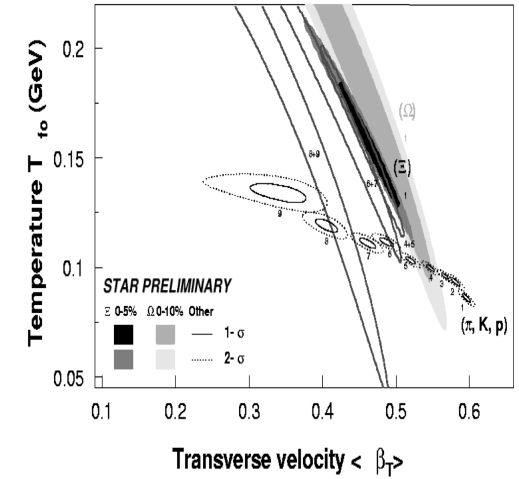
pp 7 TeV



Pb-Pb 2.76 TeV



Au-Au 200 GeV



L1

π : 0.5-1.15 GeV/c;
 K : 0.2-1.25 GeV/c;
 p : 0.3-2.30 GeV/c
 Λ : ≤ 2.75 GeV/c;
 Ξ : ≤ 3.25 GeV/c ;
 Ω : ≤ 3.0 GeV/c

L2

π : 0.5-1.35 GeV/c;
 K : 0.2-1.65 GeV/c;
 p : 0.3-2.45 GeV/c
 Λ : ≤ 2.50 GeV/c;
 Ξ : ≤ 2.70 GeV/c ;
 Ω : ≤ 3.40 GeV/c

M.Estienne, STAR Coll. arXiv:nucl-ex/0411034

Outlook

- *larger statistics => multi-differential analysis*
 - *very good PID as low as possible in p_T*
 - *charged particle multiplicity*
 - *event-shape*
 - *different ranges in $\Delta\eta$ and $\Delta\Phi$ relative to $L(T)P$*
- *Core-corona interplay in A-A and pp - plays an important role in understanding the origin of different experimentally evidenced trends*
- *pp as high as possible in charged particle multiplicity*
- *Understanding the similarities and differences between pp and A-A at high f_g^{in}*
- *lower mass A-A collisions ?*